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# Constructive interaction of $d^0$ ferromagnet with superconductor

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Recently,  $d^0$  ferromagnetism in closed shell oxides containing virtually no magnetic ions has been the subject of a number of theoretical and experimental investigations as a part of the effort to develop suitable materials for spintronic devices [1]. Although its true physical origin still remains to be solved, some surface and/or grain-boundary related defects in the nanostructures are believed to be responsible for  $d^0$  ferromagnetism [2].

In addition, although not well recognized,  $d^0$  ferromagnetism will provide an interesting experimental system for studying superconductor/ferromagnet (S/F) proximity effect. Previously, a variety of S/F heterostructures have been used to investigate the interplay between S and F order parameters [3]. However, the design and fabrication of S/F heterostructures are quite challenging because the magnetic moments in the conventional ferromagnetic materials are usually quite strong and can easily destroy the superconducting state. However, Anderson and Suhl [4] suggested long ago that a weak ferromagnetism should not destroy the superconducting state; rather, superconductivity could survive in a ferromagnetic background provided that the magnetic direction is varied on a scale smaller than the superconducting coherence length, resulting in an inhomogeneous domain-like structure called the cryptoferromagnetic state. Thus, the inherently weak and surface-derived nature of  $d^0$  ferromagnetism makes it an excellent candidate for realizing the cryptoferromagnetic state and the related domain wall effect. In this work, we hence investigate the magnetic properties of  $\text{MgB}_2/\text{MgO}$  composite, in which  $\text{MgB}_2$  and  $\text{MgO}$  phases are responsible for superconductivity and  $d^0$  ferromagnetism, respectively.

The superconductor/ $d^0$  ferromagnet composite was synthesized using solid phase reaction between Mg and  $\text{B}_2\text{O}_3$  [5]. Figure 1(a) shows the temperature dependent irreversibility of the zero-field-cooling (ZFC) and field-cooling (FC) magnetization ( $M$ ) curves measured at 50 Oe. In addition to the superconducting transition of the bulk  $\text{MgB}_2$  at 39 K, one notices two additional characteristic features at  $\sim 60$  K and  $\sim 120$  K in the ZFC magnetization curve, suggesting the existence of certain magnetic transitions at these temperatures. To highlight possible changes in the magnetization in these temperature regions, we took a difference between the  $M(H)$  loops measured at different temperatures [see Fig. 1(b)]. The difference  $M(H, T) - M(H, 125 \text{ K})$  ( $T = 90, 100, 110 \text{ K}$ ) tends to reveal a hysteresis loop characteristic of superconducting materials with weak bulk pinning. It is hence most probable that the observed drop in  $M$  at  $\sim 120$  K in the ZFC curve shows the appearance of the superconducting phase at that temperature because of the constructive interaction between  $d^0$  ferromagnet and superconductor. Furthermore, the sharp peak at  $\sim 60$  K seen in the ZFC curve is indicative of the existence of a spin-glass-like transition, which most likely results from the rearrangement of ferromagnetic domain caused, directly or indirectly, by the superconducting order occurring at temperatures below  $\sim 120$  K. We consider that the extremely high critical temperature of  $\sim 120$  K obtained in this work results from the inherently weak and inhomogeneous (domain-like) nature of defect related  $d^0$  ferromagnetism, which may induce a local increase of  $T_c$  in the vicinity of a domain wall.

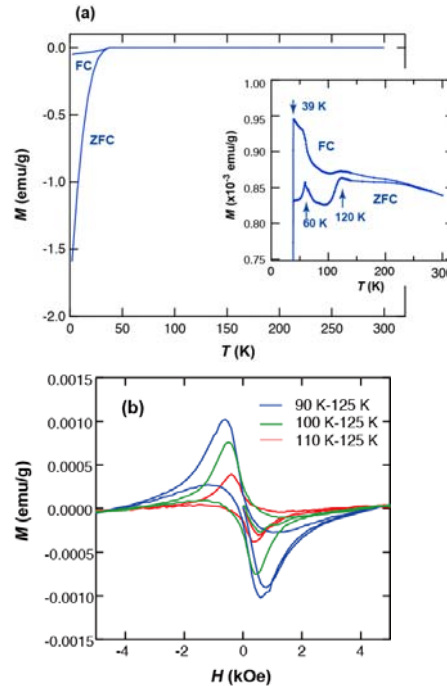


Fig. 1 (a)  $M(T)$  curves of the sample measured at 50 Oe. The inset shows an expanded plot in the temperature region above 39 K. (b) Difference between the  $M(H)$  loops measured at different temperatures indicated.

**References** [1] A. Zunger, S. Lany, and H. Raebiger, *Physics* **3**, 53 (2010). [2] T. Uchino and T. Yoko, *Phys. Rev. B* **87**, 144414 (2013). [3] A. I. Buzdin, *Rev. Mod. Phys.* **77**, 935 (2005). [4] P. W. Anderson and H. Suhl, *Phys. Rev.* **116**, 898 (1959). [5] T. Uenaka and T. Uchino, *Phys. Rev. B* **83**, 195108 (2011).