## Magnetization reversal induced by spin-orbit torque in a nanoscale Ta/CoFeB/MgO dot

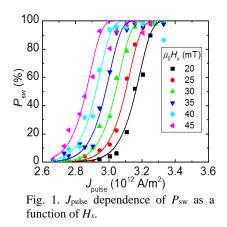
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Spin-orbit torque (SOT) induced magnetization reversal attracts increasing attention these years, because it offers a new magnetization switching scheme for spintronics devices [1]. For the development of SOT-spintronics devices, it is of particular importance to reveal factors in determining the threshold current density ( $J_{th}$ ). According to our previous work [2],  $J_{th}$  in devices with the size down to 30 nm, to which a single-domain picture is expected to be applicable, is about 3-4 times smaller than that calculated from a macrospin model [3], suggesting some factors reducing  $J_{th}$  are missing. One possible missing factor is the field-like component  $\tau_{FL}$  of SOT, which is not considered in the model studied so far. Here, we investigate the switching probability of a 40-nm Ta/CoFeB/MgO dot fabricated on a Hall-bar, and compare the result with the macrospin model including  $\tau_{FL}$ .

Figure 1 shows the switching probability  $P_{sw}$  as a function of the pulse current density  $J_{pulse}$  under different external magnetic fields  $H_x$  applied along the current direction. From the fitting with a macrospin model without  $\tau_{FL},$  the spin Hall angle  $\theta_{SH}$  and the effective anisotropy field  $\mu_0 H_{\kappa}^{\text{eff}}$  are determined to be 0.21 and 0.5 T, respectively. These values are deviated from  $\theta_{\rm SH} = 0.03$  and  $\mu_0 H_K^{\text{eff}} = 0.23 \text{ T}$  determined from other methods [4]. We simulate SOT switching by a macrospin model with the Landau–Lifshitz–Gilbert equation including  $\tau_{FL}$ , and compare the



simulated results with the experimental ones. We find that the model without  $\tau_{FL}$  overestimates both  $\mu_0 H_K^{\text{eff}}$ and  $\theta_{SH}$ , and the degree of discrepancy depends on the magnitude of effective damping constant  $\alpha_{\text{eff}}$ . Assuming  $\alpha_{\text{eff}} \approx 0.15$ , the experimentally obtained  $J_{\text{th}}$  can be well reproduced by the simulation with  $\theta_{SH} = 0.03$  and  $\mu_0 H_K^{\text{eff}} = 0.23$  T, implying that  $\tau_{FL}$  is one of the decisive factors for  $J_{\text{th}}$ .

This work was supported by R&D project for ICT Key Technology of MEXT, R&D Subsidiary Program of METI, ImPACT Program of CSTI, JSPS KAKENHI Grant No. 15J04691, and IIARE of Tohoku Univ.

[1] L. Q. Liu *et al.*, Science **336**, 555 (2012).
[2] C. Zhang *et al.*, JSAP Spring 2015, E12P-D11-5.
[3] K.-S. Lee *et al.*, APL **104**, 072413 (2014).
[4] C. Zhang *et al.*, APL. **103**, 262407 (2013).