## Quantum-mechanical effects in atomically thin MoS<sub>2</sub> FET

○Nan Fang and Kosuke Nagashio Department of Materials Engineering, The University of Tokyo E-mail: nan@adam.t.u-tokyo.ac.jp

## 1. Introduction

*C-V* measurement is the powerful tool to gain further insight into the device physics of 2D materials. We have successfully measured quantum capacitance ( $C_Q$ ) in monolayer MoS<sub>2</sub> FET<sup>[1]</sup>.  $C_Q$  is originally derived from the finite DOS of a 2D electron gas and Fermi distribution. With increasing channel thickness, the electron distribution in the finite channel thickness should contribute to the total capacitance ( $C_{\text{total}}$ ). Here, we discuss thickness dependence of quantum-mechanical effects in MoS<sub>2</sub> FET.

## 2. Results & Discussion

ALD-Al<sub>2</sub>O<sub>3</sub> (10 nm) top gated MoS<sub>2</sub> FET with different channel thickness are fabricated through the mechanical exfoliation on the quartz substrate. **Fig. 1** shows the experimental  $C_{\text{total}}$ - $V_{\text{TG}}$  curves at 1 MHz of monolayer, trilayer, 10-nm and 84-nm thick MoS<sub>2</sub>, respectively.

2.1 Accumulation region Two distinct features are observed by increasing the MoS2 thickness. Monolayer C-V shows higher saturation capacitance and sharper change than thicker MoS<sub>2</sub> case. To understand these behaviors, the quantum-mechanical model was considered by solving Poisson equation and Schrödinger equation self-consistently. Calculated square modulus of wave functions  $|\psi|^2$  for lowest subband with different MoS2 thickness are shown in Fig. 2. At thick MoS<sub>2</sub> (>tetralayer), the centroid of wave function is confined in the triangle potential formed in the channel by electrical field, which leads to the negligible thickness dependence. While at MoS<sub>2</sub> (<tetralayer) thinner than the triangle potential, the centroid of wave function is confined in the potential wall formed by the atomic channel thickness itself. This quantum-mechanical effect forces electron distribution to locate away from Al2O3/MoS2 interface. This reduces  $C_{\text{total}}$  at strong accumulation region because  $C_{\text{total}}$  is expressed as  $1/C_{\text{ox}} + l/\varepsilon_{\text{MoS2}}$ , where *l* is the distance between centroid wave function and interface. Calculated  $C_{\text{total}}$  at strong accumulation as a function of MoS<sub>2</sub> thickness is shown in **Fig. 3**. When EOT = 1 nm, the decrease of  $C_{\text{total}}$  is as high as 20 % in multilayer case. The use of monolayer MoS<sub>2</sub> is highly demanded in order to increase accumulation  $C_{\text{total}}$  especially at low EOT case.

<u>2.2 Depletion region</u> Since the parasitic capacitance is totally removed by using the quartz substrate, the minimum capacitance plateau observed for 84-nm thick MoS<sub>2</sub> results from the contribution of depletion capacitance ( $C_D$ ) with the maximum depletion width ( $W_{Dm}$ ).  $W_{Dm}$  is determined to be ~48-55 nm by thickness dependence measurement<sup>[2]</sup>. As a result, the undepleted MoS<sub>2</sub> layer will always remain, which results in residual conductance and low  $I_{ON}/I_{OFF}$  in *I-V*.  $N_D$  is determined to be 2~3×10<sup>17</sup> cm<sup>-3</sup> from  $W_{Dm}$ , which indicates the intrinsically-existed S vacancy concentration.

 $W_{\rm Dm}$  for various 2D materials are summarized here as a function of  $N_{\rm D}$  ( $N_{\rm A}$ ), as shown in **Fig. 4**. At high  $N_{\rm D}$  ( $N_{\rm A}$ ) region (>10<sup>19</sup> cm<sup>-3</sup>),  $W_{\rm Dm}$  will decrease substantially, resulting in a small thickness window for "fully controlled band transport". Full control of channel will be lost when the 2D thickness become greater than  $W_{\rm Dm}$ . This explains why well-controlled FETs with high  $I_{\rm ON}/I_{\rm OFF}$  are difficult to achieve in recent heavily doped 2D materials such as PtS<sub>2</sub>, PtSe<sub>2</sub>, SnS, and SnSe. From the above analysis, well controlled doping approaches on 2D crystals are in great demand for positive control of FET performance.

Acknowledgements This research was partly supported by the JSPS Core-to-Core Program, A. Advanced Research Networks, & JSPS KAKENHI, Japan. **References** [1] N. Fang *et al.*, J. Phys. D 2018, **51**, 065110. [2] N. Fang *et al.*, ACS Appl. Mater. Interfaces 2018, **10**, 32355.





lator Mos

Fig. 1  $C_Q$ - $V_{TG}$  curves.

Fig. 2 Calculated  $|\psi|^2$  of lowest subband .



Fig. 3 Calculated accu-

mulation  $C_{\text{total}}$  of MoS<sub>2</sub>.

C<sub>ox</sub>@ EOT = 1 nm



**Fig.4** Thickness scaling rule of 2D-FETs. 2D in green are studied here.

3.4

3.2

3.0

0.42

2.8 [he]