

Improved Response of Uni-Traveling-Carrier Photodiodes by Carrier Injection

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1. Introduction

The uni-traveling-carrier photodiode (UTC-PD), which has a p-type photo-absorption layer and a widegap electron-collection layer, has been proposed as a device to provide a higher saturation output.¹⁾ Because of the high electron velocity in the collection layer, this device is capable of operating at high current densities while still keeping its ultrafast response. The fabricated InP/InGaAs UTC-PD showed an output voltage as high as 1.9 V with a 3-dB bandwidth (f_{3db}) of 115 GHz.²⁾ The experimental results also indicated that electron transport in a p-type photo-absorption layer is of prime importance in maximizing the operating speed.

In this paper, we demonstrate that the response time of the UTC-PD's dramatically changes with the carrier injection level and the background hole density in the photo-absorption layer. The increase of f_{3db} is more pronounced in a structure with a lower hole density. This enhancement is attributed to the acceleration of electrons by the self-induced electric field associated with the hole current.

2. Experiments

The UTC-PD structures studied in this paper were grown on semi-insulating InP substrates by MOCVD. As shown in Fig. 1,

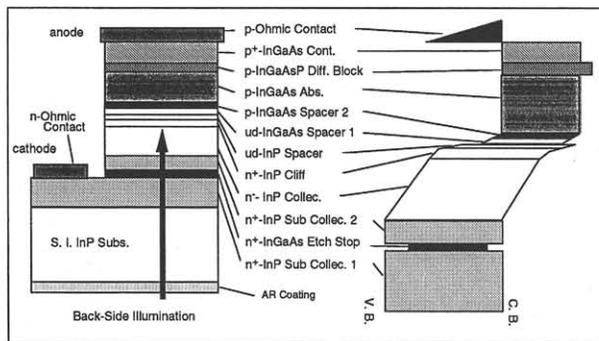


Fig. 1 Schematic cross-section and band diagram of the UTC-PD.

they consisted of an n^+ -InP/ $\text{In}_{0.53}\text{Ga}_{0.47}\text{As}$ /InP subcollector layer, followed by a 208-nm InP collector layer ($\text{Si}: 4 \times 10^{16} \text{ cm}^{-3}$), a 10-nm InP cliff layer ($\text{Si}: 2.5 - 3 \times 10^{18} \text{ cm}^{-3}$), a 2-nm undoped (ud)-InP spacer layer, a 2-nm ud- $\text{In}_{0.53}\text{Ga}_{0.47}\text{As}$ first spacer layer, a 10-nm $\text{In}_{0.53}\text{Ga}_{0.47}\text{As}$ second spacer layer ($\text{C}: 2.5 \times 10^{18} \text{ cm}^{-3}$), a 210-nm $\text{In}_{0.53}\text{Ga}_{0.47}\text{As}$ photo-absorption layer, a 15-nm p^+ - $\text{In}_{0.6}\text{Ga}_{0.4}\text{As}_{0.85}\text{P}_{0.15}$ barrier layer, and a 60-nm p^+ - $\text{In}_{0.53}\text{Ga}_{0.47}\text{As}$ cap layer. Three structures with different hole densities in the photo-absorption layer, $p = 2.5 \times 10^{17}$, 1.0×10^{18} , and $2.5 \times 10^{18} \text{ cm}^{-3}$ were prepared. The second spacer layer was inserted to keep the photo-absorption layer undepleted under operating

conditions. The UTC-PD's were fabricated using conventional photolithography, wet chemical etching, and lift-off metallization techniques. After the devices were fabricated, back-side of the InP substrates was polished and an anti-reflection coat was finally deposited for back-side illumination. The external quantum efficiency was around 22 % for all three structures.

High-speed characterizations were made for UTC-PD's with the anode electrode diameter of 5.6 μm . The device was placed at the center of a 4-mm-long 50- Ω coplanar waveguide transmission line. The signal and ground lines were respectively connected to the anode and cathode electrodes of the device, where the load resistance (R_{load}) of the UTC-PD under test was equivalent to 25 Ω . The responses against the 790-fs width laser pulse at a 1.55- μm wavelength and a repetition rate of 100 MHz were measured by using an electro-optic (EO)-sampling technique³⁾ with an external CdTe probe tip. By calibrating the EO-signal intensity using a 50-GHz digital sampling oscilloscope, the UTC-PD output peak voltages (V_{out} 's) were derived. From the Fourier transform of the measured photoresponses, the f_{3db} of the devices was evaluated.

3. Results and Discussion

Figure 2 shows the laser pulse response of the UTC-PD with $p = 1.0 \times 10^{18} \text{ cm}^{-3}$ when biased at -1.5 V. The observed V_{out} increases as the laser power increases, and it reached 1 V when the input laser power (P_{in}) was 0.99 $\mu\text{J/pulse}$. In this condition, the full width at its half maximum and the f_{3db} of the UTC-PD response were 5.1 ps and 70 GHz, respectively.

In Fig. 3, V_{out} was plotted as a function of P_{in} for three different absorption layer doping levels. Here, the bias voltage

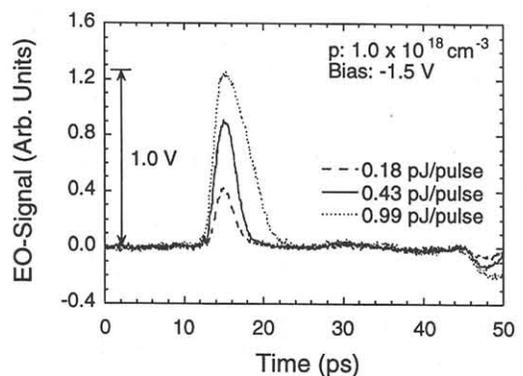


Fig. 2 Photoreponse of the UTC-PD.

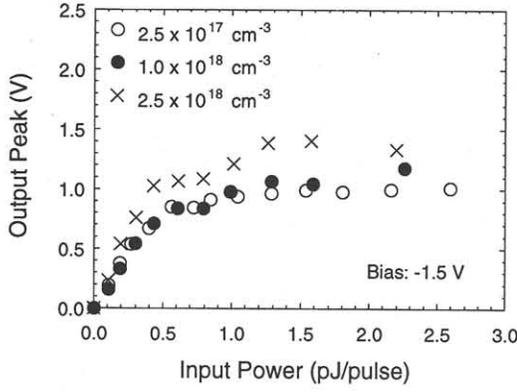


Fig. 3 Output peak voltage versus input laser power.

applied to the device was also -1.5 V. This figure shows that the $V_{out} - P_{in}$ characteristic is divided into two regions: the linear region where P_{in} is small and V_{out} is proportional to P_{in} ; and the saturation region where V_{out} scarcely changes with the increasing P_{in} . The saturation V_{out} is 1 V for the UTC-PD with $p = 2.5 \times 10^{17} \text{ cm}^{-3}$, and it increases slightly with the doping level.

The relation between f_{3dB} and V_{out} is shown in Fig. 4 for three types of UTC-PD's. The observed f_{3dB} increases first with V_{out} up to a maximum. The highest $f_{3dB} = 150 \text{ GHz}$ was obtained for the device with $p = 2.5 \times 10^{17} \text{ cm}^{-3}$. After reaching its maximum, f_{3dB} decreases with V_{out} . This f_{3dB} enhancement is more prominent for a lower doping level. It was confirmed that f_{3dB} determined by the CR-time constant was estimated to be 360 GHz for these devices at the bias voltage of less than -1.0 V. So, the difference in V_{out} dependence of f_{3dB} is attributed to the difference in carrier traveling time in the devices.

When the photo-absorption layer doping level decreases, the minority electron mobility is expected to increase. This can reduce the carrier transit time in the absorption layer and thus increase f_{3dB} . This effect, however, does not explain the enhancement of f_{3dB} with the increasing V_{out} . As the photo-excited electrons are flowing toward the cathode, a hole current is

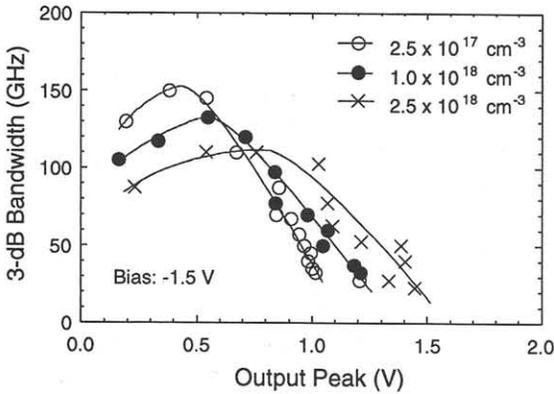


Fig. 4 3-dB bandwidth versus output peak voltage.

induced in the photo-absorption layer. Ignoring the hole diffusion component, the hole current density J_h is approximated as

$$J_h = qp_0\mu_h E, \quad (1)$$

where q , p_0 , μ_h , and E are elemental charge, the hole mobility ($\sim 150 \text{ cm}^2/\text{Vs}$), the background hole density in the photo-absorption layer, and the electric field, respectively. Equation (1) suggests that a self-induced electric field E , which simultaneously acts on electrons, is higher for a lower p_0 . It can effectively accelerate the photo-excited electrons toward the cathode electrode. Assuming that the electron current is dominated by the diffusion current in the photo-absorption layer, the potential difference over the photo-absorption layer at the output peak due to the electric field E

$$\Delta\phi_{eff} = \int E dx$$

is related to V_{out} as

$$\Delta\phi_{eff} \approx \frac{V_{out} W_A}{2q\mu_h p R_{load} S}, \quad (2)$$

where, W_A and S are the photo-absorption layer thickness and the device active area, respectively. From eq. (2), $\Delta\phi_{eff}$ is respectively estimated to be 60, 15, and 6 meV for $p = 2.5 \times 10^{17}$, 1.0×10^{18} , and $2.5 \times 10^{18} \text{ cm}^{-3}$ at $V_{out} = 0.4 \text{ V}$. As predicted in the previous paper,¹⁾ a static potential gradient of even 50 meV can considerably improve f_{3dB} . Although the self-induced electric field appears instantaneously, the f_{3dB} enhancement observed can basically be attributed to a similar mechanism.

4. Conclusion

The photoresponse of InP/InGaAs UTC-PD was systematically investigated by changing the photo-absorption layer doping level from 2.5×10^{17} to $2.5 \times 10^{18} \text{ cm}^{-3}$. Fabricated devices generated the saturation output voltage of up to 1 V at a bias voltage of -1.5 V. From the behavior of 3-dB bandwidths vs. output voltage, we found the enhanced bandwidth as a result of the self-induced electric field in the photo-absorption layer.

Acknowledgments

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