# Comparison of Sub-Bandgap Impact Ionization in Deep-Sub-Micron Conventional and Lateral Asymmetrical Channel nMOSFETs

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# 1. Introduction

Impact ionization in silicon nMOSFETs for drain voltages  $(V_D)$  well below the bandgap voltage of silicon has received widespread attention [1,2,3]. Substrate currents  $(I_{SUB})$  for drain voltages down to 0.6V [1] and floating body effects in SOI devices down to 0.8V [2] were reported. This would imply that the impact ionization induced operational and reliability issues in nMOSFETs will continue to deca-nano meter device generations.

Based on Monte Carlo simulations it was suggested that various modes of electron-electron interactions resulting in the high energy tail (HET) of the electron energy distribution are responsible for some electrons to have more energy than that gained from the lateral electric field (E<sub>LAT</sub>) [3,4]. An anomalous increase of the gate voltage at which the ISUB peaks (V<sub>Gpeak</sub>) which can not be explained by HET presented. We have also compared the sub-bandgap impact ionization in CONventional (CON) and Lateral Asymmetrical Channel (LAC) nMOSFETs of channel length 100nm. An enhancement of the increase in V<sub>Gneak</sub> is found in the LAC devices. Based on the results presented we propose quantization of inversion layer as an additional energy gain mechanism for the electrons.

## 2. Experimental

The MOSFETs used in this study had a channel length of 100nm and gate oxide thickness of 3.6nm. Both CON and LAC MOSFETs were fabricated on the same wafer for fair comparison. The fabrication procedure of the devices are described in detail elsewhere [5]. Fig. 1 shows the simulated channel doping profiles of the CON and LAC MOSFETs used in this study and reveals the nonuniform channel doping for the LAC devices. Fig. 2 shows the output characteristics of the CON and LAC devices.

# 3. Results and Discussions

The I<sub>SUB</sub>-V<sub>G</sub> characteristics of the devices were measured for V<sub>D</sub> down to 0.85V. Fig. 3 shows the I<sub>SUB</sub>-V<sub>G</sub> plots for CON and LAC for the lowest V<sub>D</sub> investigated and the LAC shows much lower I<sub>SUB</sub> than the CON. In fig.4 the V<sub>Gpeak</sub> is compared. The V<sub>T</sub> is subtracted from V<sub>Gpeak</sub> to account for the difference in V<sub>T</sub> of the two devices. The V<sub>Gpeak</sub> shows expected linear behavior for V<sub>D</sub> above 1.5V. Below 1.5V, the V<sub>Gpeak</sub> is found to deviate from this and start increasing as the V<sub>D</sub> is decreased further. For LAC the increase in V<sub>Gpeak</sub> is much more pronounced for low V<sub>D</sub> than the CON. Fig. 5 compares the ratio I<sub>SUB</sub>/I<sub>D</sub> at I<sub>SUBpeak</sub> for the two devices. Also shown is the ratio of I<sub>SUB</sub>/I<sub>D</sub> falls off more rapidly as V<sub>D</sub> is lowered as compared to the high V<sub>D</sub> regime. This fall-off is less rapid for the LAC than the CON as seen from the I<sub>SUBpeak</sub> ratio.

Fig. 6 shows the  $I_{SUB}/I_D$  versus  $1/V_D$  plot for both CON and LAC where it is seen that the data deviates from the

predictions of the lucky electron model for low V<sub>D</sub>. It was shown that Auger recombination can be an additional energy gain mechanism [6]. Fig. 7 shows the correlation between  $I_{SUB}$  and  $I_D^2 I_{SUB}$  [6]. Although a great part of the data for  $V_D=1.5V$  support the Auger recombination as an additional energy gain mechanism, for low V<sub>D</sub> such a correlation is not found. The high energy tail theories in the present form [3,4] can not explain the V<sub>G</sub> dependence presented in fig. 4.

Figures 8, 9 and 10 show the simulated  $E_{LAT}$ , transverse field ( $E_{TRA}$ ) and electron concentration respectively for  $V_D$ =0.9V and  $V_{Gpeak}$ - $V_T$  for CON (0.48V) and LAC (0.87V). For the LAC the  $E_{LAT}$  is smaller than that for CON. The  $E_{TRA}$ near the drain becomes smaller as the  $V_G$  is increased. The effect of the increase in the positive  $E_{TRA}$  is to pull up the electrons more to the interface as shown in the fig. 10.

#### 4. Model

Based on the results presented above we propose inversion layer quantization as an additional energy gain mechanism for the electrons. The concept is illustrated in fig. 11. The electrons in the quantized inversion layer has energy higher than the conduction band edge, E<sub>C</sub>. But the electrons can relax to E<sub>C</sub> near the drain junction where there is no quantized layer. As can be appreciated from fig. 11, the electrons in the quantized case effectively have a higher energy than otherwise [7]. As the V<sub>D</sub> is reduced the energy gain from ELAT decreases, which calls for stronger inversion. The enhanced V<sub>Gpeak</sub> increase for LAC can be due to the lower ELAT and hence the energy deficit to cause impact ionization is more. The electron concentration at the surface would increase for the CON devices too for larger V<sub>G</sub> as illustrated in fig. 10. The electron mobility near the drain end is much smaller for the CON than the LAC because of an order of magnitude higher channel doping, fig. 1 and larger field, figures 8 and 9, near the drain [8]. This makes the secondary energy gain mechanisms less efficient in the case of CON as is evident from the more rapid fall of ISUBneak for V<sub>D</sub> below 1.3V for CON, fig. 5, right y-axis.

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Fig. 1 Comparison of doping distribution in CON and LAC MOSFETs.



Fig. 4 Dependence of the gate voltage at which  $I_{SUB}$  peaks ( $V_{Gpeak}$ ) on the  $V_D$  for CON and LAC MOSFETs.



Fig. 7 I<sub>SUB</sub> versus  $I_D^2 I_{SUB}$  showing a power law relationship for a great part of the data for VD=1.5V for both CON and LAC suggesting Auger recombination as a possible energy gain mechanism. For low V<sub>D</sub> this dependence is very weak.

![](_page_1_Figure_6.jpeg)

**Fig. 10** Distribution of electron concentration near the drain junction for LAC and CON devices for the bias conditions described in fig. 8.

![](_page_1_Figure_8.jpeg)

**Fig. 2**  $I_D$ - $V_D$  characteristics of the CON and LAC MOSFETs. LAC shows better saturation and higher current drive.

![](_page_1_Figure_10.jpeg)

Fig. 5 Dependence of the  $I_{SUB}/I_D$  ratio on  $V_D$ . The ratio decreases rapidly below  $V_D$  of 1.3V. The ratio of  $I_{SUB}$  of LAC to that of CON is also shown.

![](_page_1_Figure_12.jpeg)

Fig. 8 The lateral electric field distribution for both the devices corresponding to the  $V_{Gpeak}$  conditions for the CON and LAC shown in fig. 4.  $V_{GT} = 0.48V$  correspond to  $V_{GT}$  peak for CON and 0.87V for the LAC.

![](_page_1_Figure_14.jpeg)

Fig. 3  $I_{SUB}$ - $V_G$  plots for CON and LAC MOSFETs measured for  $V_D$  well below band-gap voltage of silicon.

![](_page_1_Figure_16.jpeg)

Fig. 6 The  $I_{SUB}/I_D$  ratio versus  $1/V_D$  plot. At low  $V_D$  the data deviates from the straight line indicating that LEM is inadequate to explain  $I_{SUB}$  at low  $V_D$ .

![](_page_1_Figure_18.jpeg)

**Fig. 9** The transverse field distribution for LAC and CON devices for the bias conditions described in fig. 8.

![](_page_1_Figure_20.jpeg)

Fig. 11 Inversion layer quantization as an energy gain mechanism. In the quantized inversion layer the electrons occupy the sub-bands (marked 0 and 1). Near the drain junction the electron layer is not quantized. In this region the electrons can relax to the bottom of the conduction band resulting in an effective energy gain corresponding to the energy of the sub-band above  $E_c$ .