

E-1-3

The Effect of Emitter Size Scaling on  $f_{max}$  in InP/InGaAs HBTs

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1. Introduction

InP-based HBTs are promising candidates for such applications as high speed ICs for 40Gbps optical communication systems. While high  $f_T$  can be readily obtained in these HBTs because of superior electron transport of InP-related materials,  $f_{max}$  depends rather on processing technique as well as device size and is thus less predictable compared to  $f_T$ . For instance, while several reports have shown that  $f_{max}$  improves with decreasing emitter width [1], it has also been reported that  $f_T$ , which is one of the important factors determining  $f_{max}$ , deteriorates when emitter size is down scaled beyond certain limit [2]. Thus in order to achieve high  $f_{max}$ , it is crucial to understand the device scaling of  $f_{max}$  in correlation with the behavior  $f_T$ .

In this paper, we investigate the emitter size dependence of  $f_{max}$  of single and double heterojunction HBTs (SHBTs and DHBTs) with different  $f_T$ . With optimum device design, a state-of-the-art  $f_{max}$  of 313 GHz was obtained. Critical design consideration in improving the  $f_{max}$  of InP-based HBT will be discussed.

Table I Epitaxial layer structures of SHBTs and DHBTs

| Layer         | Material               | Doping  | Thickness(nm) |      |
|---------------|------------------------|---------|---------------|------|
|               |                        |         | SHBT          | DHBT |
| Emitter       | n <sup>+</sup> -InGaAs | Si:3e19 | 100           | ←    |
|               | n <sup>+</sup> -InP    | Si:2e19 | 20            | ←    |
|               | n-InP                  | Si:3e17 | 80            | ←    |
| Base          | p <sup>+</sup> -InGaAs | C:4e19  | 50            | ←    |
| Collector     | i-InGaAs               | Si:1e16 | 300           | 100  |
|               | n <sup>+</sup> -InP    | Si:1e18 | -             | 20   |
|               | i-InP                  | Si:1e16 | -             | 180  |
| Sub-collector | n <sup>+</sup> -InGaAs | Si:1e19 | 400           | ←    |

2. Device fabrication

The epitaxial layers of HBTs studied in this work were grown by MOCVD. Table I shows epitaxial layer structures of SHBT as well as DHBT with ballistic launcher structure in the collector region [3]. The former type of HBT has typical  $f_T$  of 160 GHz, whereas in the latter case  $f_T$  exceeds 200 GHz. To reduce base resistance, the base layer was heavily doped to  $4 \times 10^{19} \text{ cm}^{-3}$ . Prior to formation of base

contact, the wafers were annealed at 500 C to remove hydrogen out from the carbon-doped base layers [4]. As shown in Fig.1, the HBTs were fabricated using self-aligned process. The emitter width,  $W_E$ , ranged from 0.2  $\mu\text{m}$  to 2.6  $\mu\text{m}$ , whereas the width of each base finger,  $W_B/2$ , was fixed at 0.5  $\mu\text{m}$ .

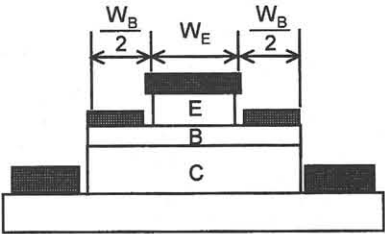


Fig. 1. Schematic cross section of fabricated HBT

3. Experimental Results

The RF performance of the HBTs was characterized by on-wafer measurement of S-parameters from 0.1 to 40 GHz.  $f_{max}$  was obtained by extrapolating Mason's unilateral gain (MUG) with -6 dB/Oct. slope as shown in Fig. 2.

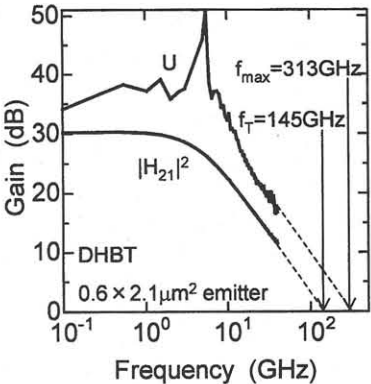


Fig.2. Frequency dependence of  $|H_{21}|$  and MUG

Fig.3 plots  $f_T$  and  $f_{max}$  of SHBTs and DHBTs as a function of  $W_E$ . The emitter length was fixed at 2.1  $\mu\text{m}$ . As shown in Fig.3, both type of HBTs show similar  $W_E$  dependence of  $f_{max}$ . As long as  $W_E$  is greater than 1  $\mu\text{m}$ ,  $f_{max}$  is inversely proportional to square root of  $W_E$ . This  $W_E$  dependence of  $f_{max}$  indicates ideal scaling relationship for small scale HBTs. However, as  $W_E$  is decreased to below 1  $\mu\text{m}$ , the variation of  $f_{max}$  with  $W_E$  deviates from ideal

scaling and turns to decreasing trend. Qualitatively, this behavior of  $f_{\max}$  can be understood by degradation of  $f_T$ , also shown in Fig.3.

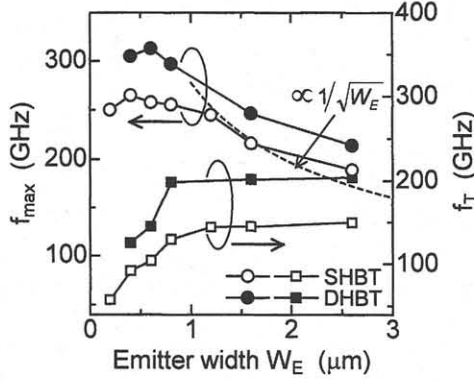


Fig.3. Emitter width dependence of  $f_{\max}$

Regarding maximum  $f_{\max}$ , while the  $f_{\max}$  of SHBT peaked at 265 GHz, in DHBT peak  $f_{\max}$  of 313 GHz was obtained (Fig. 2). Note that the optimum  $W_E$  is far below the size at which  $f_T$  starts to drop. This indicates that in order to optimize device size for high  $f_{\max}$  a specific design guideline taking all relevant parameters into account is needed.

#### 4. Discussion

In order to gain insight into the factors responsible for the  $f_{\max}$  behavior,  $f_T$  was plotted as a function of  $R_B C_{BC}$  as shown in Fig.4. The time constant  $R_B C_{BC}$  was estimated from  $f_T$  and  $f_{\max}$  using the common relationship between  $f_T$  and  $f_{\max}$ :

$$f_{\max} = \sqrt{\frac{f_T}{8\pi R_B C_{BC}}} \quad (1)$$

This plot allows us to determine which of the two parameters,  $f_T$  or  $R_B C_{BC}$ , is dominating the  $f_{\max}$ . Figure 4 indicates that when  $W_E$  is greater than 1.0 μm  $f_{\max}$  is mainly determined by  $R_B C_{BC}$ . As a result,  $f_{\max}$  increases with decreasing  $W_E$ . The maximum  $f_{\max}$  can be obtained by drawing a line tangent to the plotted curve, then evaluating  $f_{\max}$  from the line slope using the following relationship which is equivalent to eq. (2).

$$f_T = (8\pi f_{\max}^2) R_B C_{BC} \quad (2)$$

Note that  $R_B C_{BC}$  continues to decrease even after  $f_{\max}$  starts to decrease. Therefore, when  $W_E$  is excessively decreased, it is  $f_T$  rather than  $R_B C_{BC}$  that is limiting  $f_{\max}$ .

To obtain more specific design guidelines for high  $f_{\max}$  HBT, we performed a simple analysis. Taking the emitter size dependence of  $f_T$  and  $R_B C_{BC}$  into account, it can be shown that the optimum  $f_{\max}$  as well as the emitter width that determines the

optimum point are expressed as follows:

$$f_{\max}^{opt} \cong \frac{1}{4\pi\sqrt{R_B^{ext} C_{BC}^{ext}}} \left( \tau_F + 2\rho_c^E c_0 \left( 1 + \sqrt{1 + \frac{\tau_F}{\rho_c^E c_0}} \right) \right)^{-1/2} \quad (3)$$

$$W_E^{opt} \cong W_B \sqrt{\frac{\rho_c^E c_0}{\tau_F + \rho_c^E c_0}} \quad (4)$$

|                |                                                   |
|----------------|---------------------------------------------------|
| $R_B^{ext}$    | : External base resistance                        |
| $C_{BC}^{ext}$ | : External collector capacitance                  |
| $\tau_F$       | : Total carrier transit time delay                |
| $\rho_c$       | : Specific contact resistivity of emitter contact |
| $c_0$          | : Collector capacitance per unit area             |

From (3) and (4), it can be clearly understood that the optimum  $f_{\max}$  is determined by ratio  $\tau_F/\rho_c^E c_0$ . If the emitter contact is good ( $\tau_F \gg \rho_c^E c_0$ ),  $\tau_F$  is the dominant factor that determines  $f_{\max}^{opt}$ . In this case, significant increase in  $f_{\max}$  can be expected by improving  $f_T$  and by further pushing the down scaling of  $W_E$ . On the other hand, with poor contact ( $\tau_F \ll \rho_c^E c_0$ ), decreasing  $W_E$  below  $W_B$  proves to be useless and  $f_{\max}^{opt}$  is insensitive to  $\tau_F$  (and thus to  $f_T$ ). The ratio  $\tau_F/\rho_c^E c_0$  in our HBTs is about 3. Thus, for further improvement of  $f_{\max}$ , improving not only  $\tau_F$  but also emitter contact is effective.

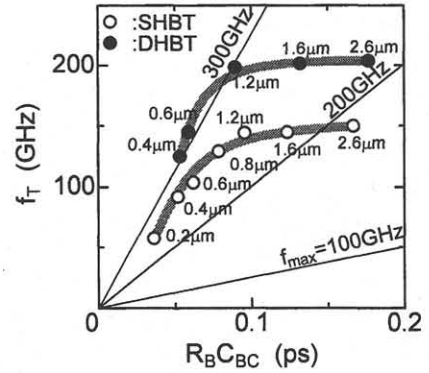


Fig.4.  $f_T$  versus  $R_B C_{BC}$  plot

#### 5. Conclusions

Emitter size scaling in InP-based HBTs for high  $f_{\max}$  was investigated. Using simple analysis, we clarified the essential key factors in improving  $f_{\max}$ . While an excellent  $f_{\max}$  of 313 GHz was obtained with our HBT, the analysis suggests that further improvement is possible by improving the emitter contact.

#### Reference

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