## Control of ferromagnetic order in selectively *p*-doped GaMnAs-based heterostructures

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While the bandgap and wavefunction engineering are limited semiconductor far to nonmagnetic SO heterostructures, we aim to broaden its use to magnetic heterostructures and to extend the degree of freedom in designing spin-related properties in semiconductors. In this paper, we use the two-dimensional hole gas (2DHG) system in selectively-doped GaAs/p-AlGaAs heterostructures together with  $\delta$ -doping of magnetic (Mn) impurities, and successfully maximize the ferromagnetic order among the Mn spins in GaAs by overlapping the wavefunction of 2DHG with the Mn  $\delta$ -doping profile.

The mainstream studies of spin-electronic materials on III-V semiconductors consist based of (i) ferromagnet/semiconductor heterostructures such as MnAs/GaAs<sup>1</sup> and (ii) magnetic alloy semiconductors such as InMnAs and GaMnAs.<sup>2-4</sup> Despite extensive studies, in the former system special techniques are required to grow multilayer heterostructures with abrupt interfaces,<sup>5</sup> and much higher Curie temperature  $T_{\rm C}$  is needed for practical application in the latter system (the highest  $T_{\rm C}$  for the past few years was 110K for (GaMn)As).<sup>6</sup> Unlike the random alloy system, here we use  $\delta$ -doping of Mn in GaAs. Inherent advantages of  $\delta$ -doping<sup>7</sup> are locally high dopant concentration and high carrier concentration, which can lead to high Curie temperature  $T_{\rm C}$ .<sup>8</sup> Another prospective advantage is easy fabrication of multilayer heterostructures containing Mn  $\delta$ -doped GaAs layers with excellent interfaces.<sup>9,10</sup>

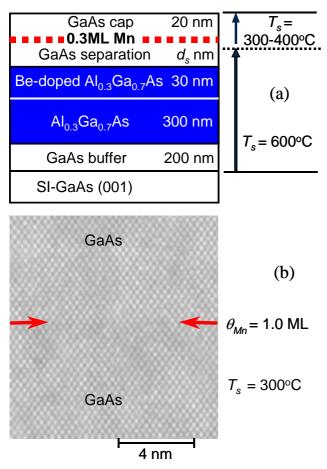
Mn  $\delta$ -doped GaAs layers were grown on semi-insulating (SI) GaAs (001) substrates by MBE at the growth temperatures  $T_s = 200-400^{\circ}$ C. Extensive structural analyses revealed that most of the Mn atoms are abruptly confined within a width of 2-3 monolayers (ML) in the zinc-blende structure as substitutional dopants, when the nominal thickness  $\theta_{Mn}$  of Mn is below 1 ML.<sup>10,11</sup> The Mn doping profiles retained abruptness even at elevated  $T_s$  up to 400°C. Although it was possible to incorporate high Mn concentration in the Mn  $\delta$ -doped GaAs layers, the hole to Mn concentration ratio  $p/\theta_{Mn}$  was very low and was not enough to realize ferromagnetic ordering.<sup>10,11</sup>

In order to obtain high hole concentration and locally high Mn concentration at the same position, we have grown 0.3 ML Mn  $\delta$ -doped GaAs/Be-doped *p*-AlGaAs heterostructures by MBE, whose structure is shown in Fig. 1 (a). These *p*-type selectively doped heterostructures (*p*-SDHS) resemble an inverted high electron mobility transistor, where holes are provided from the *p*-AlGaAs layer to the overlying GaAs layer. The thickness  $d_s$  of the undoped-GaAs separation layer was a measure to control the interaction between the Mn  $\delta$ -doped GaAs layer and the 2DHG formed at the GaAs/*p*-AlGaAs interface.

Hall loops at 40 K of 0.3 ML Mn  $\delta$ -doped GaAs layers grown at 400°C without and with p-SDHS\_( $d_s = 3$  nm) are shown in Fig. 2 (a) and (b), respectively. The hysteresis in the Hall loop of the sample with p-SDHS clearly indicates ferromagnetic order, while the ferromagnetic hysteresis is absent in the sample without p-SDHS. The temperature dependence of the sheet resistance ( $R_{\text{sheet}} - T$ ) of the samples without and with p-SDHS is plotted in Fig. 2 (c). The sample without p-SDHS shows insulating behavior due to the low hole concentration. In contrast, the sample with p-SDHS shows a local maximum of the  $R_{\text{sheet}} - T$  trace, suggesting that  $T_{\text{C}}$  is ~70 K. This value of  $T_{\text{C}}$  was confirmed by measuring Hall loops at various temperatures, where hysteresis remained open up to 70 K.

The ferromagnetic order of the samples with p-SDHS was found strongly dependent on  $d_s$ . As shown in Fig. 2 (d), the local maximum temperature of the bump in the  $R_{\text{sheet}} - T$  trace, which roughly corresponds to  $T_{\text{C}}$ , was 45 K at  $d_s = 0$  nm and 70 K at  $d_s = 3$  nm. The Hall loops showed clear ferromagnetic hysteresis at  $d_s = 0$  and 3 nm below  $T_{\rm C}$ . With further increase of  $d_{\rm s}$  to 5 and 10 nm, the bump disappeared. Hysteresis in loops was not observed at  $d_s = 5$  and 10 nm, indicating the absence of The  $d_s$  dependence of the ferromagnetic order. ferromagnetic order is explained using the valence band diagram of the heterostructure in Fig. 2 (e). We think that the degree of the overlap of the 2DHG wavefunction and the Mn  $\delta$ -doping profile directly affects the ferromagnetic ordering and  $T_{\rm C}$  of the heterostructures.  $T_{\rm C}$  was highest (70 K) at  $d_s = 3$  nm, because the overlap was maximum. Further increase of  $d_s$  to 5 and 10 nm decreased the overlap of the 2DHG wavefunction and the Mn  $\delta$ -doped profile, thus weakened the ferromagnetic order.

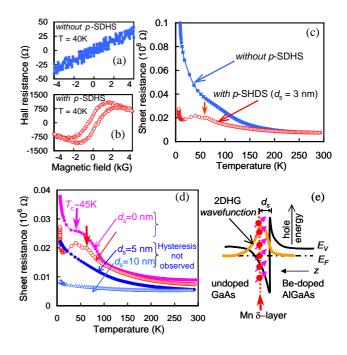
In order to completely suppress the surface segregation of Mn, and to obtain locally higher Mn concentration in a ideally sharp  $\delta$ -doped profile,  $T_s$  of the Mn  $\delta$ -doped GaAs layer in the SDHS was lowered from 400°C (at which surface segregation of around 30% Mn dopants was observed in the SIMS depth profile) to 300°C (at which no



**Fig. 1.** (a) Sample structure of Mn δ-doped GaAs with *p*-type selectively doped heterostructures (*p*-SDHS). The GaAs separation layer thickness  $d_s$  was 0 - 10 nm. Holes are supplied from the Be-doped *p*-type Al<sub>0.3</sub>Ga<sub>0.7</sub>As layer to the Mn δ-doped GaAs layer in the SDHS. (b) High-resolution TEM lattice image of the Mn δ-doped GaAs layer with  $\theta_{Mn} = 1.0$  ML grown at  $T_s = 300^{\circ}$ C.

segregation of Mn was detected).<sup>10,11</sup> The sample examined here was a 0.3 ML Mn  $\delta$ -doped GaAs ( $T_s = 300^{\circ}$ C)/Be-doped *p*-AlGaAs heterostructure with  $d_s = 0$  nm, as shown in Fig. 1 (a). Low-temperature (LT) annealing was carried out in a nitrogen atmosphere for 15 minutes with various annealing temperatures  $T_a = 280 - 340^{\circ}$ C. It was estimated by the Curie-Weiss fitting that  $T_C$  is as high as 172 K for the sample annealed at  $T_a = 300^{\circ}$ C. Note that  $T_C$  was dramatically increased from 70 K (grown at  $T_s =$ 400°C; as-grown) to 112 K ( $T_s = 300^{\circ}$ C; as-grown), and 172 K ( $T_s = 300^{\circ}$ C; LT-annealed).

In summary, we have shown that the controlled overlap of the wavefunction of the 2DHG and the Mn  $\delta$ -doping profile in GaAs can lead to ferromagnetic ordering. The highest  $T_{\rm C}$  of the ferromagnetic heterostructures prepared with suitable growth conditions and low-temperature annealing was 172 K. This  $T_{\rm C}$  value is the highest among the reported values in III-V (GaAs, InAs) magnetic semiconductors. Furthermore, we will show that the ferromagnetic order in the *p*-SDSH can be controlled by gate electric field and light irradiation.



**Fig. 2.** (a) & (b) Hall loops of 0.3 ML Mn  $\delta$ -doped GaAs layers grown at 400°C *without* and *with p*-SDHS ( $d_s = 3$  nm), measured at 40 K. (c)  $R_{\text{sheet}} - T$  traces of the samples *without* and *with p*-SDHS, respectively. (d)  $R_{\text{sheet}} - T$  traces of 0.3 ML Mn  $\delta$ -doped GaAs samples *with p*-SDHS for  $d_s = 0, 3, 5$  and 10 nm. (e) Schematic diagram of the valence band profile of the *p*-SDHS, the 2DHG wavefunction, and Mn dopants.  $E_V$  and  $E_F$  are the valence band top and the Fermi energy, respectively. *z* is the growth direction.

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