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## Leakage mechanism of ultrathin SiON gate dielectric

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### Abstract

Leakage mechanism of ultrathin gate SiON dielectric is studied considering dangling-bonds in the dielectric. The calculated and measured  $J_G$ -EOT characteristics are compared. Additionally, larger flat band shift in pMOS is studied using the first principles method.

Keywords: N≡Si<sub>3</sub>, SiON, dangling-bond, TAT, and thin EOT

### Introduction

Ultrathin SiON film is a candidate for next gate dielectric, because the improvement of fabrication process has decreased equivalent oxide thickness (EOT) and gate leak [1], suppressed boron penetration [2] and flat band shift ( $\Delta V_{FB}$ ) [3], and removed interfacial roughness [1]. However, still unclear is why  $\Delta V_{FB}$  is larger in pMOS, as shown in Fig. 1. We will solve this problem and then study how dangling-bond (DB) affects the gate leak through the SiON dielectric.

### Alloy Structure of SiON

Fig. 2 shows the ternary component diagram of SiON that is composed of  $[(\text{SiO}_2)_{1-X}(\text{Si}_3\text{N}_4)_X]_{1-Y}\text{Si}_Y$  including interfacial transition (IFT) layers, in which energy gap (EG), dielectric constant (K), valence band affinity (VBA) are changed from in Si to at the edges of the dielectric [4-6]. Throughout the dielectric, physical quantities (Q) are dependent of Si-N bond rate (R) [7, 8] but not of [N] [9] as follows: (Q in SiON) =  $(1-R) \times Q_{OX} + R \times Q_{SiN}$ , where  $R=3X/(1+2X)$  and  $X=0.75[N]/(1-[N])$ . The profiles of [N] and [O] are measured by AR-XPS method. We can then extract R, EG, K, VBA, IFT layer widths ( $W_{IFT}$ ), and tunnel masses of electrons and holes ( $m_e$  and  $m_h$ ) from the measured [N]-profile by using the values for Q in SiO<sub>2</sub> and Si<sub>3</sub>N<sub>4</sub> ( $Q_{OX}$  and  $Q_{SiN}$ , respectively) in Table 1. Fig. 3(a)-(d) depicts the results of a measured sample.

### Analysis of dangling-bonds and flat band shift

In Fig. 4, the first principles calculation shows that N-atom incorporated in Si (100) causes Si dangling bonds (Si-DB). In Fig. 5, the upper insertion depicts Si-DB and 2-fold N dangling-bond (N-DB) in the Si energy band [3]. A combined trap distribution of these DBs has a broad tail in the density-of states (DOS) diagram. According to [10], the center of the combined distribution ( $E_T$ ) is assumed to be at the level lower from the Si mid-gap by 0.4eV at the interface. The half-value width (0.1eV [10]) of the distribution can be extended with  $\Delta E_T(z) = 0.1eV \times EG(z)/1.12eV$ , where z is a position in space through the dielectric and EG(z) is shown in Fig. 3 for example. In Fig. 6, the positive charge is contributed by upper tail above the Fermi level ( $E_F$ ) at each bias condition, and then causes the  $V_{FB}$  shift and tunnel barrier modulation. Since the upper tail is composed mainly of Si-DBs, the  $V_{FB}$  shift and the barrier modulation can be ascribed to the Si-DBs.

Dividing the film into grids ( $i=1, 2\dots M$ ) and integrating the Poisson equation, we have  $K_i \varepsilon_0 \Delta V_i / \Delta T_i = -q \Delta N_i$ , where  $\varepsilon_0$  is the vacuum permittivity, q is the elementary charge, and  $\Delta T_i$  is the width of i-th grid layer.  $\Delta N_i$  and  $\Delta V_i$  are the number of trapped charges of i-th grid layer and potential drop due to  $\Delta N_i$ , respectively.  $K_i$  is the dielectric constant of i-th grid layer. Integrating these grid layers, we have  $\Delta V_{FB} = \sum \Delta V_i = -(q/M/\varepsilon_0) \gamma T_{phys}$ , where  $\gamma = \sum \Delta N_i / (3.6R_i - 3.9)$  and  $T_{phys}$  is a physical thickness of dielectric film. Since only a portion of Si-N bonds gives trap site, we cannot directly extract  $\Delta N_i$  from the measured R-profile. Then, we assume that the trap density ( $N_T$ ) is equal to  $Y_{DB} \times R$ , where  $Y_{DB}$  is the dangling-bond (trap site) yield from Si-N bonds. In Fig. 7,  $Y_{DB}$  is extracted to reproduce  $V_{FB}$  in measured CV characteristics and the resultant values are summarized in Table 2. Note that  $Y_{DB}$  in pMOS is

2-times larger than in nMOS, which is equivalent to the behavior of  $\Delta V_{FB}$  in Fig. 1. In pMOS, larger  $\Delta V_{FB}$  is therefore ascribed to larger  $Y_{DB}$ .

### Calculation of gate leakage current

Fig. 8 depicts the impact of barrier modulation estimated as tunnel barrier change from that of  $Y_{DB}=0$  (no positive charge). This modulation increases the barrier for holes and lowers that for electrons, and is also dependent of bias condition, as schemed in Fig. 6.

On the other hand,  $m_h$  has been unknown, while  $m_e$  was extracted to be  $0.85m_0$  with including the IFT layers [4], where  $m_0$  is the rest electron mass. To extract  $m_h$ , we compare calculated and measured currents of samples in which the [N]-profile is uniform [12], as shown in Fig. 9. There exist cross points in currents between electrons and holes, since the increase of [N] substantially reduces VBA, that is, tunnel barrier for holes, from 4.49eV to 1.9eV (See Table 1). It is found that the calculated cross points are agreed with the measured one while  $m_e=m_h=0.85m_0$ .

Trap-assisted tunneling (TAT) current is calculated by integrating the product of: (1) capture cross-section of traps ( $\sigma$ ), (2)  $N_T(z)$ , and (3) 2-step tunneling probabilities of electrons and holes [11]. The integration is carried out with space (z) and energy (E). The gate current density ( $J_G$ ) is then calculated by adding this to the direct tunneling current of electrons and holes that aren't captured by any traps. In Fig. 10, TAT disappears while the capture radius is in atomic scale.

### Results

Taking into account  $m_e=m_h=0.85m_0$ , the tunnel barrier modulation, the measured [N]-profile, the extracted  $Y_{DB}$ , and regarding  $\sigma$  as a fitting parameter, we compare the CV-JV curves between experiments and theory. An excellent agreement is obtained when  $\sigma=0$ , as shown in Fig. 11. The leakage mechanism is therefore direct tunneling affected by positive charges trapped by Si-DBs. In Fig. 12, we compare the calculated and measured  $J_G$ -EOT characteristics. To suppress the gate leak, nMOS requires larger  $Y_{DB}$  because  $T_{phys}$  increased for a given EOT overcomes the barrier lowering (shown in Figs. 6 and 8) partially increasing the electrons current. In opposite, pMOS requires smaller  $Y_{DB}$  because  $T_{phys}$  increased for a given EOT overcomes the reduction of the barrier increasing for holes (shown in Figs. 6 and 8). Additionally, B atoms of p<sup>+</sup>poly-Si can generate B-N-B cluster in the dielectric. Since this cluster is mobile even in β-Si<sub>3</sub>N<sub>4</sub> in which the cluster is more stable than in SiON [18], the extrinsic Si-DB is left after the cluster has gone, as shown in Fig. 13. Therefore, pMOS has the extrinsic Si-DBs as well as the intrinsic Si-DBs generated before fabricating the gate poly-Si, while nMOS has only the intrinsic Si-DBs.

### Conclusion

If  $Y_{DB}$  is controlled, smaller EOT and smaller gate leak can be realized. We are required to obtain larger  $Y_{DB}$  in nMOS and smaller  $Y_{DB}$  in pMOS. Boron penetration from p<sup>+</sup>poly-Si can generate the extrinsic  $Y_{DB}$ , which results in larger  $\Delta V_{FB}$  in pMOS. The leakage mechanism is the direct tunneling affected by Si-DBs.

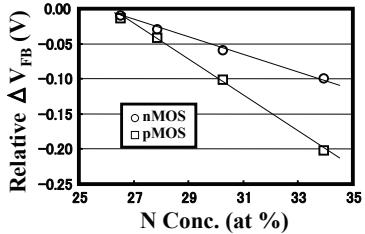
### References

- [1] Matsushita, VLSI04, 172. [2] Yang, IEDM99, 245. [3] Matsushita, IEDM05, 847. [4] Watanabe, SSDM04, 732, and to appear in TED53, issue6, 2006. [5] Muller, Nature399, 758, 1999. [6] Giustino, PRL91, 267601, 2003. [7] Lukovsky, SSDM96, 356. [8] Yasuda, SSDM01, 486. [9] Guo, EDL19, 207, 1998 [10] Yasuda, SSDM01, 202. [11] Yasuda, SSDM93, 847. [12] Muraoka, JAP94, 2038, 2003. [13] Miyazaki, Appl. Surf. Sci. 113/114, 585, 1997. [14] Alay, JAP81, 1606, 1997. [15] Miyazaki, Appl. Surf. Sci. 216, 252, 2003. [16] "Physics of Semiconductor Devices", Sze, Wiley, 1981. [17] Matsushita, PHD Thesis, U. Nagoya, 2001. [18] Kato, to appear in ICPS06 Vienna, July, 2006.

**Table 1 Used Variable**

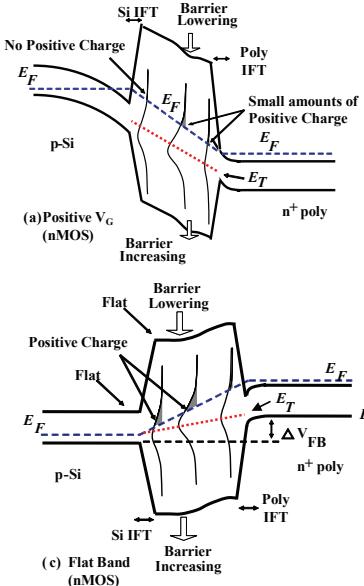
Quantity (Q)	$Q_{\text{OX}}$	$Q_{\text{SiN}}$	$Q_{\text{Si}}$
K	3.9 [16]	7.5 [16]	11.7
EG (eV)	8.95 [13]	5.4 [15]	1.12
VBA (eV)	4.49 [14]	1.9 [15]	0
$m_e (m_0)$	0.85 [4]	0.85 (*)	—
$m_h (m_0)$	0.85 (*)	0.85 (*)	—
W <sub>IFT</sub> (nm)	0.4 [4-6]	0.5 [17]	—

(\*) Present results (Fig. 9)



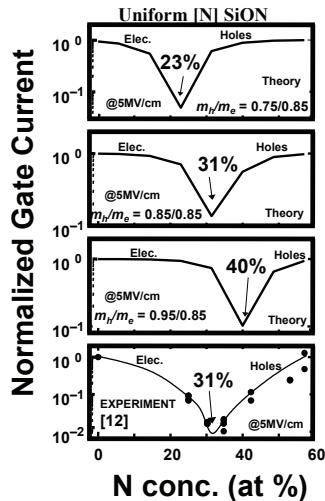
**Fig. 1 Relative measured  $V_{\text{FB}}$  shift**

Derivative in pMOS is 2-time larger than in nMOS.



**Fig. 6 Schemes for Trapped Positive Charge**

The modulation of EG profile shown in Fig. 3 was omitted for ease of drawing in this figure. If this modulation is considered, the barriers for electrons and holes are further lowered and partially lowered, respectively. Of course, such an impact of EG modulation on the barrier modulation is considered in the calculation.

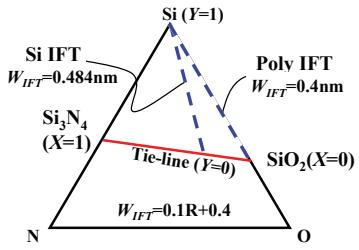


**Fig. 9 Tunnel mass for holes**

**Table 2 Extracted parameters**

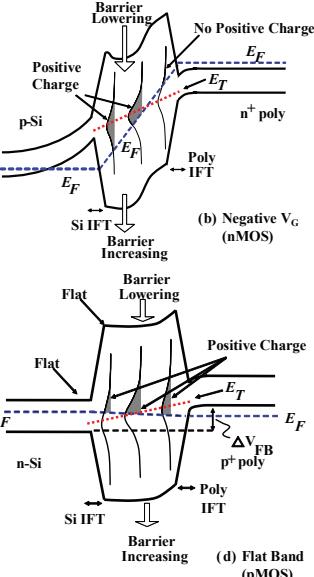
Sample	EOT	0.78 nm	0.89 nm	1.0 nm
Type	pMOS	nMOS	nMOS	

$Y_{\text{DB}}$  in pMOS is 2-times larger than that in nMOS.



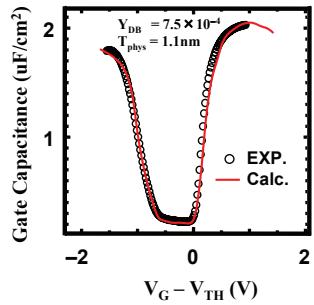
**Fig. 2 Ternary component diagram of SiON**

$W_{\text{IFT}}$  is dependent of R. Tie line is on Y=0 (away from IFT layers).

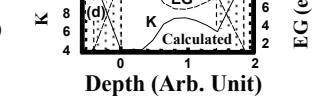


**Fig. 3 Profiles of SiON dielectric of a measured sample**

EG and K are changed with [N] and [O] through the dielectric film including the IFT layers. Physical thickness ( $T_{\text{phys}}$ ) is a distance between the centers of IFT layers [4].

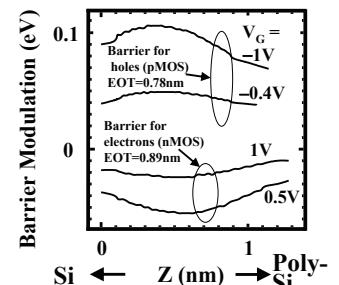


**Fig. 4 First principles result of Si-DB due to incorporating N in Si(100)**



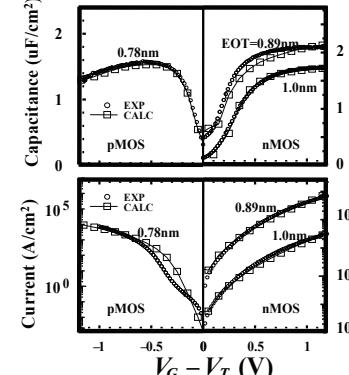
**Fig. 5 Model for trap levels in SiON**

$\text{VBA}_{\text{OX}}$ ,  $\text{VBA}_{\text{SiN}}$ ,  $\text{EG}_{\text{OX}}$ ,  $\text{EG}_{\text{SiN}}$  and  $\text{EG}_{\text{Si}}$  are from the literature (Table 1).

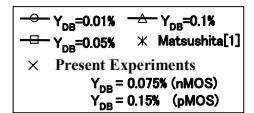


**Fig. 7 CV-fitting to extract  $Y_{\text{DB}}$**

$Y_{\text{DB}}$  is determined to reproduce  $V_{\text{FB}}$ .

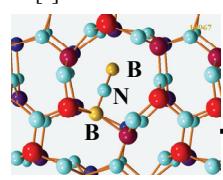


**Fig. 8 Barrier modulation from  $Y_{\text{DB}}=0$**



**Fig. 9 CV-JV fitting**

Impurity density profiles measured by SIMS are used in the calculation. All the CV-JV characteristics are calculated using the Poisson-Schrödinger solver developed by adding the present DB model to that used in [4].

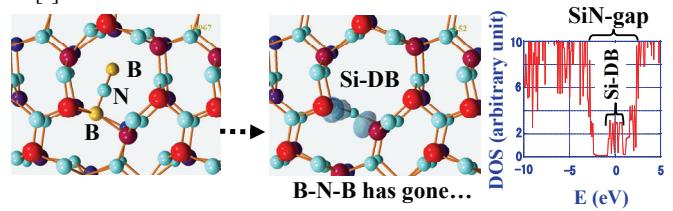


**Fig. 10 Impact of TAT on  $J_G$**

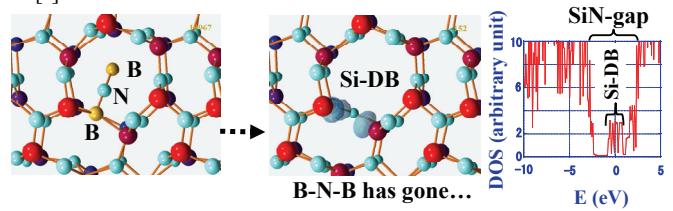
TAT of holes is considered as well as that of electrons. The barrier modulation is considered and  $m_e = m_h = 0.85m_0$  is used in the calculation.



**Fig. 11 CV-JV fitting**



**Fig. 12 Impact of  $Y_{\text{DB}}$  on  $J_G$ -EOT**



**Fig. 13 Generation of extrinsic Si-DB from B-N-B cluster in  $\beta\text{-Si}_3\text{N}_4$**