

# Effect of Insertion Layers on Schottky Barrier Height of Fe/*n*-type Ge Junctions

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## Abstract

We investigated the effect of thin insertion layers of GeO<sub>2</sub> and MgO on the Schottky barrier height ( $\phi_{\text{SBH}}$ ) of the Fe/*n*-type Ge junctions where the Ge has a non-degenerate character. Current-voltage (*I-V*) characteristics of the Fe/GeO<sub>2</sub>/Ge devices showed a significant suppression of the  $\phi_{\text{SBH}}$  whereas the *I-V* curves hardly depended on thickness of the MgO layer. We report a large reduction of the  $\phi_{\text{SBH}}$  from 0.60 eV down to about 0.25 eV by inserting the GeO<sub>2</sub> layer.

## 1. Introduction

Technologies for electrical spin injection in practical semiconductors (SCs) using a ferromagnetic (FM) tunnel contact have recently made remarkable progress [1]. For Si and Ge, the spin injection/detection experiments have been mainly done in degenerately doping SCs. From an application point of view, however, a non-degenerate SC should be more important because it is utilized as the core part of the practical SC devices such as a channel layer of metal-oxide-semiconductor field-effect-transistor (MOSFET). For a metal/non-degenerate SC junction, one of the major obstructs for developing a practical device should be a formation of wide depletion region due to the Schottky barrier potential. This causes a high resistance-area-product (*RA*), thereby a degradation of the device performance in a high-speed operation.

Ge is a promising candidate as a channel layer material of next-generation high-speed MOSFET. The Fermi-level ( $E_{\text{F}}$ ) at a metal/Ge interface is strongly pinned at around the top of the valence band of Ge, resulting in high Schottky barrier height ( $\phi_{\text{SBH}}$ ) of 0.5 ~ 0.6 eV for *n*-type Ge regardless of the work function of the metal electrode [2]. Nishimura *et al.* has recently reported that GeO<sub>2</sub> is one of the most effective insulating materials to alleviate the  $E_{\text{F}}$ -pinning and reduce the  $\phi_{\text{SBH}}$  for the non-magnetic metal (Al and Au)/*n*-Ge junctions [3].

In this study, we demonstrate that the GeO<sub>2</sub> insertion layer is quite effective to reduce the  $\phi_{\text{SBH}}$  for the FM metal (Fe)/non-degenerate *n*-type Ge junction.

## 2. Experiments

Films were grown by molecular beam epitaxy on non-degenerate phosphorous-doped Ge(001) substrates with a carrier concentration of  $\sim 5 \times 10^{16} \text{ cm}^{-3}$  at room temperature (RT). Tunnel contacts consisting of a Au(10 nm)/Fe(3 nm)/GeO<sub>2</sub>( $d_{\text{GeO}_2}$  nm) were deposited by elec-

tron-beam evaporation at RT. As reference samples, epitaxial Fe(001)/MgO(001) tunnel contacts were also grown on the substrate at RT, where the MgO is a canonical tunnel barrier [1,4]. The junctions with several different sizes of active area (*A*) between 25×12.5 and 600×300  $\mu\text{m}^2$  were prepared with standard micro-fabrication techniques.

## 3. Results and discussion

Figures 1(a) and 1(b) show the current-voltage (*I-V*) characteristics of the devices with the GeO<sub>2</sub> and MgO insertion layers at RT. The current dramatically increases by inserting GeO<sub>2</sub> layers up to  $d_{\text{GeO}_2} = 1.0 \text{ nm}$ , indicating the alleviation of the  $E_{\text{F}}$ -pinning and a large reduction of the  $\phi_{\text{SBH}}$ . The current starts to decrease with further increasing  $d_{\text{GeO}_2}$ . In contrast to the GeO<sub>2</sub>-based device, no noticeable change was observed in the *I-V* characteristics of the Fe/MgO/Ge devices, suggesting that the  $\phi_{\text{SBH}}$  is hardly

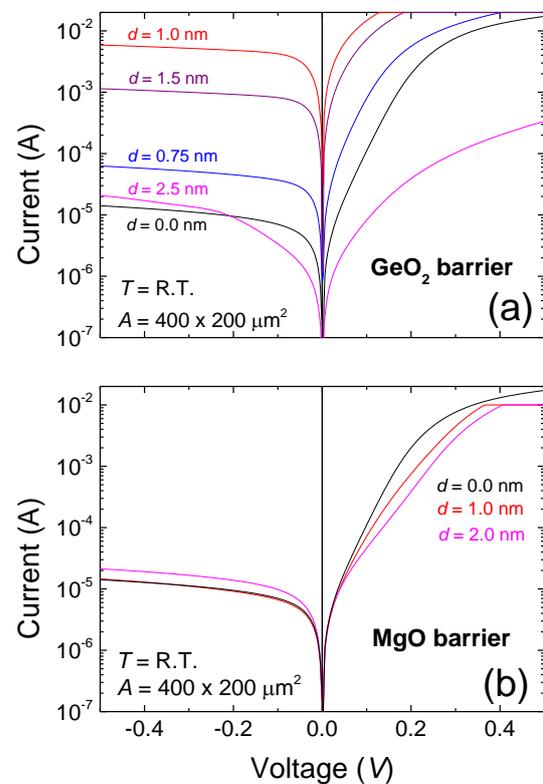


Fig. 1 *I-V* characteristics of the devices with (a) GeO<sub>2</sub> and (b) MgO insertion layers at RT, respectively.

suppressed by inserting the MgO layer. This is consistent with our previous results showing an absent of Schottky potential in the Fe/MgO/*p*-type Ge junctions [5, 6].

The *RA* as a function of  $d_{\text{GeO}}$  is given in Fig. 2, where the *RA* is obtained from a linear part of *I-V* characteristics at a low *V* range of  $\pm 10$  mV. For  $d_{\text{GeO}} \geq 1.5$  nm, a good linear relation in the  $\ln(RA)-d_{\text{GeO}}$  plot appeared, which is characteristics in a tunnel junction [7]. This indicates that tunneling is a dominant transport process and the *RA* is governed by the GeO<sub>2</sub> layer for  $d_{\text{GeO}} \geq 1.5$  nm.

For a conventional metal/SC Schottky diode, the thermionic emission is a dominant transport process. Then, the value of  $\phi_{\text{SBH}}$  can be estimated from the slope of the  $\ln(I_S/T^2)-T^{-1}$  plot (thermal-activation method) without considering the tunneling transport [7], where the  $I_S$  is the saturation current. In a case of the present Fe/GeO<sub>2</sub>/Ge devices, however, this method is no longer valid to evaluate an accurate value of the  $\phi_{\text{SBH}}$  because the tunneling is a major transport mechanism, especially for  $d_{\text{GeO}} \geq 1.5$  nm. Here, we estimate the  $\phi_{\text{SBH}}$  of the Fe/GeO<sub>2</sub>/*n*-Ge devices by taking the tunneling probability of the GeO<sub>2</sub> layer into account. For a metal/insulator/non-degenerate SC junction,  $RA \equiv (dJ/dV)^{-1}$  at low *V* is expressed by [7]

$$RA \approx \frac{\eta k}{A^* T q} \exp(\zeta^{1/2} d) \exp\left(\frac{q\phi_{\text{SBH}}}{kT}\right) \quad (1),$$

where  $\eta$  is the ideality factor,  $k$  is the Boltzmann constant,  $A^*$  is the effective Richardson constant,  $T$  is the temperature,  $q$  is the electron charge. The  $\exp(-\zeta^{1/2} d)$  represents the tunneling probability due to the insulating layer. Here,  $\zeta$  (in eV) and  $d$  (in angstrom) are the effective barrier height and thickness of the insulating layer, respectively (A constant of  $1.01 \text{ eV}^{1/2}/\text{angstrom}$  is omitted). Following the equation (1), a relative change in the  $\phi_{\text{SBH}}$  from that of the device with  $d = 0.0$  nm ( $\Delta\phi_{\text{SBH}}$ ) can be given by

$$\Delta\phi_{\text{SBH}} \approx \frac{kT}{q} \left[ \ln\left(\frac{A^*(0)}{A^*(d)} \cdot \frac{\eta(d)}{\eta(0)} \cdot \frac{RA(0)}{RA(d)}\right) + \zeta^{1/2} d \right] \quad (2).$$

The value of  $\zeta$  is estimated to be 0.35 eV from the slope of linear part in the  $\ln(RA)-d_{\text{GeO}}$  plot. We obtained the  $\phi_{\text{SBH}}$  ( $= 0.60 - \Delta\phi_{\text{SBH}}$ , where 0.60 eV is the  $\phi_{\text{SBH}}$  of Fe(001)/*n*-Ge(001) junction [8]) as a function of  $d_{\text{GeO}}$  by assuming that the  $A^*$  does not depend on  $d_{\text{GeO}}$  (Fig. 2). The  $\phi_{\text{SBH}}$  decreases and reaches to be  $\sim 0.25$  eV (corresponding to  $\Delta\phi_{\text{SBH}} \sim 0.35$  eV) for  $d_{\text{GeO}} \geq 1.5$  nm. Further reduction of the  $\phi_{\text{SBH}}$  would be expected by optimizing growth conditions, post-annealing or using lower work-function FM electrode such as Gd [9].

#### 4. Conclusion

In this study, we present a systematic investigation on the effect of the GeO<sub>2</sub> and MgO insertion layers on the  $\phi_{\text{SBH}}$  and *RA* for the Fe/non-degenerate *n*-type Ge junctions. The

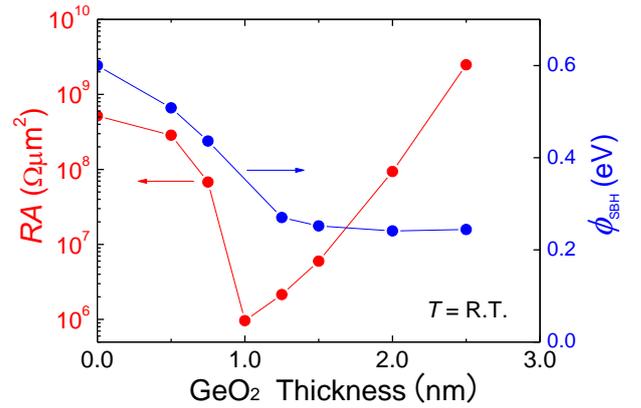


Fig. 2 *RA* and  $\phi_{\text{SBH}}$  of the Fe/GeO<sub>2</sub>/*n*-Ge devices as a function of the GeO<sub>2</sub> thickness.

data on *I-V* characteristics show a large suppression of the  $\phi_{\text{SBH}}$  by inserting the GeO<sub>2</sub> layer. A good linear relation was observed in the  $\ln(RA)-d_{\text{GeO}}$  for  $d_{\text{GeO}} \geq 1.5$  nm, indicating that the *RA* is determined by the FM tunnel contact. For the device with  $d_{\text{GeO}} \geq 1.5$  nm, the values of  $\phi_{\text{SBH}}$  were estimated to be about 0.25 eV by taking the tunneling probability of the GeO<sub>2</sub> into account. This result offers that the GeO insertion layer is a promising candidate for the tunnel barrier of the FM contact for Ge-based spintronic device.

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