# Ge Oxidation does not follow the Deal-Grove Mechanism

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## 1. Introduction

The dominance of silicon as a semiconductor is partly based on the well-behaved dielectric properties of its native oxide SiO<sub>2</sub>. Ge is a possible high mobility channel to replace Si, but its native oxide is much less well behaved. One of the benefits of Si was its well-understood oxidation mechanism that followed the Deal-Grove[1] mechanism. The oxidation of Ge somehow differs, it creates Ge/GeO<sub>2</sub> interfaces with more interface defects, molecular GeO is volatile [2] and the oxidation kinetics have a strange pressure dependence [3], but the Ge/GeO<sub>2</sub> interface can be flat.

There are three possible processes for the oxidation of Si or Ge, diffusion of  $O_2$  molecular interstitials (denoted  $O_2^*$ ) along channels in the SiO<sub>2</sub>, diffusion of interstitial lattice oxygens ( $O_i$ ), or diffusion of oxygen vacancies ( $V_O$ ) along the SiO<sub>2</sub> lattice, Fig. 1. Silicon follows the Deal-Grove model [1], in which  $O_2^*$  diffuses along the interstitial channels of the SiO<sub>2</sub> network to react exothermically with Si at the Si/SiO<sub>2</sub> interface [4]. This is possible because of the remarkably open lattice of SiO<sub>2</sub>. This is verified by the lack of  $O^{18}$  exchange with network oxygens [5]. On the other hand, O exchange with the network is seen during Ge oxidation [6,7]. This suggests that Ge oxidation follows a different mechanism. This is also seen by the different temperature dependence of Ge oxidation [7,8,2,3], Fig. 2.

#### 2. Method

Here, we use density functional supercell models to calculate defect formation and migration energies [9], as in Table 1. The total diffusion barrier is the sum of these two energies. The O chemical potential ( $pO_2$ ) changes across the film, and this greatly changes the defect formation energies.

### 3. Results.

The energy of  $O_2^*$  in the interstitial channels depends on O-O closed shell repulsions. Fig 3 shows that this energy decreases rapidly with interstitial volume V as V<sup>-3.5</sup>, and follows the same dependence in both SiO<sub>2</sub> and GeO<sub>2</sub>.The larger O bond angle in a-SiO<sub>2</sub> allows it to have large order rings of bonds, and so the interstices can expand more easily as the O<sub>2</sub>\* passes. This lowers the diffusion energy, Fig. 4. GeO<sub>2</sub> has a smaller O bond angle, which makes the channels narrower, the network stiffer, and this raises the diffusion energy to ~2.79 eV, Table 1.

The alternative path is by network interstitial oxygens,

 $O_i$  or 'peroxyls'. The energy of the peroxyl is relatively constant but depends slightly on the Si-Si separation. The diffusion transition state is the same for SiO<sub>2</sub> and GeO<sub>2</sub> and consists of a 5-fold Gs or Ge site. In GeO<sub>2</sub>, this process costs slightly less than in SiO<sub>2</sub> and it becomes the lowest energy pathway in a-GeO<sub>2</sub>, Fig. 5.

At high pO<sub>2</sub>, the interstitial oxygen of either O<sub>2</sub>\* or O<sub>i</sub> form dominates. But at low pO<sub>2</sub>, near the interface, it crosses over to diffusion by oxygen vacancies, Fig. 8. O vacancies are a costly defect in SiO<sub>2</sub>, even at the Si/SiO<sub>2</sub> interface, because it involves breaking two Si-O bonds, Table 1. But the vacancies have a low cost at the Ge/GeO<sub>2</sub> interface at this equilibrium, and this allows off-stoichiometry. Nevertheless, the diffusion barrier is high, and so overall interstitials dominate except very close to the interface. In Table 1, we compare the energy barriers of the dominant process in bold to the experimental activation energies of oxidation.

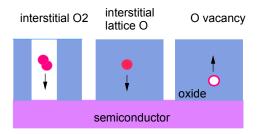


Fig 1. Diffusion processes, molecular O interstital  $(O_2^*)$ , network O interstitial, O vacancy.

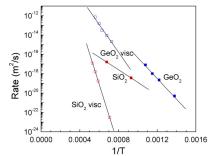


Fig. 2. Experimental oxidation rates of Si and Ge and kinematic viscosities of  $SiO_2$ ,  $GeO_2$ 

#### References

- 1. B E Deal, A S Grove, JAP 36 3770 (1965)
- 2. S K Wang, A Toriumi, Jpn J App Phys 50 04DA01 (2011)
- 3. C H Lee et al, APX 5 114001 (2012)
- 4. A Bongiorno, A Pasquarello, PRL 93 086102 (2004)
- 5. F Rochet, S Rigo, et al, Adv Phys 35 237 (1986)
- 6. S R M daSilva, I Baumvol, et al, APL 100 191907 (2012)
- 7. G Urbain, Geochim. Cosmochim. Acta 46 1061 (1982)

- S K Sharma, D Virgo, I Kushiro, J Non-Cryst Solids 33 235 (1979)
- 9. H Li, J Robertson, IEEE-SISC (San Diego, Dec 2016) p7.2
- W Xu, T Nishimura, T Yajima, A Toriumi, JSAP meeting (March 2017), 16p-413-4

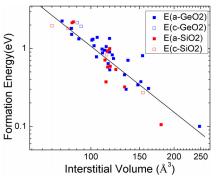


Fig. 3. O2\* defect formation energy vs. relaxed interstitial volume.

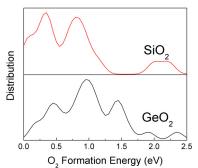


Fig. 4. Probability distribution of O2\* formation energy in SiO2 and GeO2.

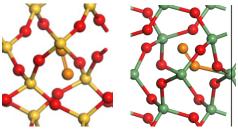


Fig. 5.  $O2^*$  migration path in SiO2 and GeO2. green = Ge,

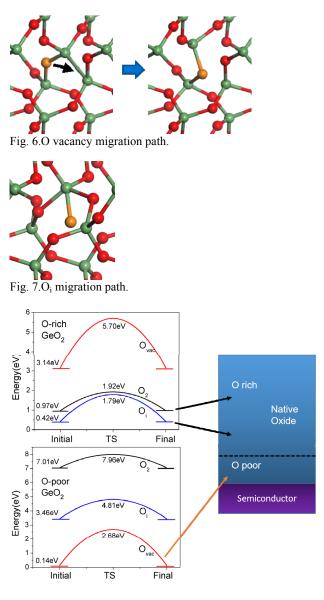


Fig. 8.Summary of diffusion mechanisms and energy barriers in different parts of the GeO2 film.

	a-GeO <sub>2</sub>			a-SiO <sub>2</sub>		
Bond angle	130 <sup>0</sup>			144-150 <sup>0</sup>		
Bulk heat of formation/O	3.02			4.85		
Heat of formation (exp)	3.06			4.78		
	Formation energy (eV)	Migration barrier (eV)	Diffusion energy (eV)	Formation energy (eV)	Migration barrier (eV)	Diffusion energy (eV)
O <sub>2</sub> interstitial / O <sub>2</sub>	0.97	1.92	2.89	0.7	0.7	1.40
O interstitial <sup>§</sup>	0.42	1.37	1.79	1.40	1.39	2.79
O vacancy (O-rich)	3.16	2.54	5.7	6.05	4.28	10.33
O vacancy (O-poor)	0.14	2.54	2.68	1.20	4.28	5.48
Experiment [1,3]			2.00			1.23

Table 1. Calculated defect formation energies and diffusion energies for a-SiO2 and a-GeO2, compared to experiment.