

A Simple Efficient Method of Nanofilm-on-Bulk-Substrate Thermal Conductivity Measurement Using Raman Thermometry

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Abstract

Using Raman thermometry, we established a simple efficient method of thermal conductivity (k) measurement for light-absorbing nanofilms on bulk non-absorbing substrates with low k . We determine film k comparing experimental and theoretical laser-induced film heating calculated assuming diffusive phonon transport (DPT). We show that the method works for films with thickness h larger than the phonon mean free path Λ , even for low k materials like nanocrystalline Si. For $h < \Lambda$, despite ballistic phonon transport contribution, this approach works when the in-plane DPT dominates, e.g. in Si films on quartz with $h \geq 60$ nm. Our method is simpler and more efficient than any other existing NF k measurement method, a dozen of NFs being examined in one hour.

1. Introduction

Measurement of the thermal conductivity (k) of nanofilms (NFs) and, especially, Si-based NFs is an important task for science and industry. Bulk Si was never used in thermoelectric applications due to its high $k \sim 150$ W/m/K and, therefore, low thermoelectric figure of merit $ZT < 0.01$ at room temperature. A reduction of k in Si NFs and nanowires (NWs) due to phonon boundary scattering and other effects can improve their performance. Further k reduction can be achieved in nanocrystalline Si-based NFs, which are composites of Si and Ni silicide nanocrystals (Ni-Si NC). Due to both phonon scattering at the grain boundaries and planar defects, the films have much lower k (3 – 7.7 W/m/K) and higher ZT (> 0.1) than bulk Si at room temperature [1]. In this work, we examine k in Si NFs and composite NC NFs of Si and Fe silicide (Fe-Si NC) using Raman thermometry. Recently, Raman scattering was utilized to measure k in nano-membranes (NMs) [2] and suspended NWs [3]. Situation is more complex for NFs on bulk substrates playing a role of the heat drain. Raman spectroscopic measurement of the k of such NFs is still a challenge. In this work, we solve this problem for the case of absorbing NF on bulk non-absorbing low- k substrate and measure k of Si NFs and composite Fe-Si NC NFs on quartz.

2. Theoretical and Experimental Methods

Let us consider an absorbing layer with complex refractive index $N_1 = \sqrt{\varepsilon_1} = n_1 + i\kappa_1$ on an infinite substrate

with the refractive index $N_2 = \sqrt{\varepsilon_2} = n_2$. Laser light hits the absorbing layer from media with the refractive index $N_0 = \sqrt{\varepsilon_0} = n_0$. Density of the light power flux is

$$S = (P/\pi r_0^2) \exp(-r^2/r_0^2) \quad (1)$$

where P is the laser power while r_0 is the radius of the focused laser light beam (Fig. 1).

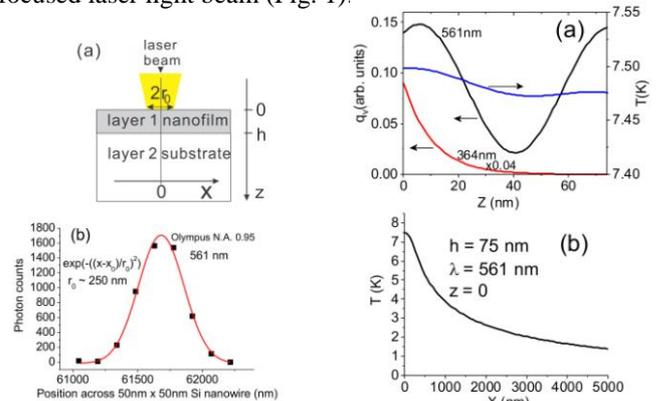


Fig. 1. a) Laser illumination of absorbing NF on transparent substrate; b) Measured light intensity distribution in the focused 561 nm laser beam (black squares) and its Gaussian fitting (red curve)

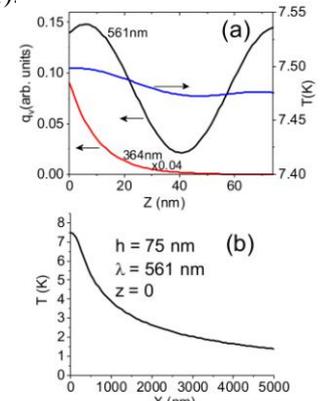


Fig. 2. a) Volumetric heat density $q_v(z)$ at $x=0$ in 75 nm thick SOQ by the 561 nm (black curve) and 364 nm (red curve) light absorption; the temperature field $T(z)$ (blue curve) for the 561 nm light with $P = 1$ mW; b) Temperature field $T(x)$ at $z=0$ for 75 nm thick SOQ.

The light is propagating from the semi-space 0 ($z < 0$), hitting the absorbing layer 1 ($0 \leq z \leq h$) and propagating further into semi-space 2 ($z > h$). Volume density of heat generated by the light absorption can be expressed as

$$q_v = \varepsilon_1'' |E_1|^2 \omega / (8\pi) \quad (2)$$

where ω is the light frequency, ε_1'' is the imaginary part of dielectric constant while E_1 is the electric field of light in the layer 1. We calculate E_1 first. Then we solve stationary heat transfer problem for layered medium with temperatures T_1 and T_2 thermal conductivities k_1 and k_2 in layers 1 and 2, respectively. The governing equations are

$$\left(\frac{\partial^2}{\partial z^2} + \frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial}{\partial r} \right) T_1 = -\frac{q_v}{k_1} \quad (0 \leq z \leq h) \quad (3)$$

$$\left(\frac{\partial^2}{\partial z^2} + \frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial}{\partial r} \right) T_2 = 0 \quad (z \geq h) \quad (4)$$

$$\text{For } z=0 \quad \frac{\partial T_1}{\partial z'}=0; \text{ for } z=h \quad T_1=T_2 \quad k_1 \frac{\partial T_1}{\partial z'}=k_2 \frac{\partial T_2}{\partial z'}$$

$$\text{At } r \rightarrow \infty \quad T_1=T_2=0; \text{ at } z \rightarrow \infty \quad T_2=0$$

After Hankel transform, we obtain a system of equations that we solve using Green's function method. Finally, we calculate an averaged laser-induced heating ΔT_{calc} of the illuminated area of layer 1.

$$\Delta T_{calc} = \int_{V_1} T_1(\vec{x}) q_V(\vec{x}) d^3\vec{x} / \int q_V(\vec{x}) d^3\vec{x} \quad (5)$$

Raman measurement was done using a Nanofinder 30 confocal Raman system (Tokyo Instruments Inc.) equipped with a 561 nm wavelength laser. Lens with 100× magnification and 0.95 numerical aperture was used to focus laser light onto spot with r_0 , ~ 250 nm. ΔT was determined from the laser-induced Raman downshift $\Delta\omega$.

For testing our method, we measured k of Si on quartz (SOQ) and Si on sapphire (SOS) NFs. Their thinning was done by thermal oxidation with subsequent etching in HF. Composite FeSi₂₀-2P NFs were fabricated using phase separation from amorphous Fe-Si alloy films that had a composition ratio of Si/Fe ~ 20 similar to Ni-Si [1]. The amorphous films were deposited on fused silica substrates by DC magnetron sputtering of a target. The sputter deposition was carried out in pure Ar gas at a sputtering pressure of 3.0×10^{-3} Torr. The thickness of the deposited films was controlled by deposition time [1]

3. Results

Figure 2a shows calculated heat volume density $q_V(z)$ at $x=0$ in 75 nm thick SOQ induced by the 561 nm (black curve) and 364 nm (red curve) light absorption and the temperature field $T(z)$ (blue curve) for the 561 nm light with $P = 1$ mW. As temperature field $T(x)$ at $z=0$ (Fig. 2b shows, the heat travels a long distance in the x-direction before being absorbed by the substrate.

Figure 3 shows the power dependencies of the laser-induced heating and their linear fittings for 75 nm thick SOQ and SOS. From the slopes of the dependencies, we find $\Delta T/P \sim 6.7$ (K/mW) for SOQ and $\Delta T/P \sim 1.9$ (K/mW) for SOS. Determination of the corresponding k_1 values is graphically shown in Fig. 4. Straight horizontal lines represent experimental dimensionless $P/\Delta T/\lambda/k_2$ values for the 75 nm thick SOQ and SOS while curves show $P/\Delta T/\lambda/k_2$ values calculated for a wide range of k_1/k_2 ratios. The k_1 can be found from the intersection between theoretical curve and corresponding horizontal line. The $k_1 \sim 65.5$ W/m/K obtained for SOQ appears to be in good agreement with literature while that of SOS ~ 22 W/m/K does not.

Figure 5 shows the theoretical and experimental Si NF thickness dependencies of k : The theoretical calculations were made by assuming in-plane DPT and frequency-dependent boundary scattering [4] Our data (open circles) for SOQ with $h \geq 60$ nm appeared to be in good agreement with theoretical and experimental in-plane Si NF k values from Ref. [4]. However, for SOQ with $h < 60$ nm, k_1 rapidly drops with a decrease in h , probably, indicating

significance of the cross-plane ballistic phonon transport. In this case like in case of SOS, our theoretical model needs modification.

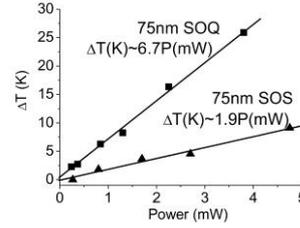


Fig. 3. 75 nm thick SOQ and SOS $\Delta T(P)$ dependencies and their linear fittings.

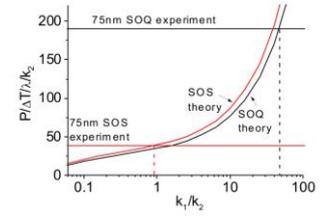


Fig. 4. Calculated dimensionless $P/\Delta T/\lambda/k_2$ parameters vs. k_1/k_2 for 75 nm thick SOQ (black curve) and SOS (red curve). Horizontal lines represent experimental $P/\Delta T/\lambda/k_2$ values. Vertical dashed lines show the determined k_1/k_2 for both cases.

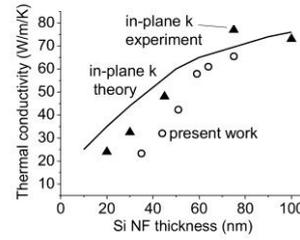


Fig. 5. Thickness dependence of Si NF thermal conductivity k : 1) theoretical and experimental results for in-plane k [4] (black curve and triangles, respectively), and 2) results obtained in this work (open circles).

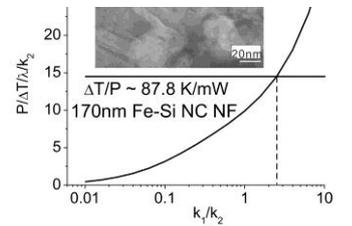


Fig. 6. Dependence of calculated dimensionless parameter $P/\Delta T/\lambda/k_2$ on k_1/k_2 for 170 nm thick Fe-Si NC NF and experimentally determined value of the parameter (horizontal line). Vertical dashed line shows the determined k_1/k_2 . Inset shows TEM image of Fe-Si NC NF film cross-section.

Now let us consider the case with $h > \lambda$ and $k_1/k_2 \sim 1$. Figure 6 shows the k_1 determination for 170 nm thick Fe-Si NC NF. The inset shows a TEM image of Fe-Si NC NF cross-section with few tens nanometer grain sizes. The λ in NC NF is even lower than the grain size. The NF displays very high $\Delta T/P \sim 87.8$ K/mW and exhibits low $k_1 \sim 3.5$ W/m/K in agreement with the value obtained using the thermal reflectance method. Despite $k_1/k_2 \sim 2.5$, the solution is stable (Fig. 6) due to sufficiently large h .

4. Conclusions

We established a Raman spectroscopic k measurement method for absorbing NFs on bulk transparent substrates and applied it for Si and FeSi₂₀-2P NFs on SiO₂.

References

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