

# Tunnel Nitride Passivated Contacts for Silicon Solar Cells Formed by Cat-CVD

Yuli Wen, Huynh Thi Cam Tu and Keisuke Ohdaira

JAIST

1-1 Asahidai, Nomi, Ishikawa 923-1292, Japan  
Phone: +81-761-51-1565 E-mail: s2020040@jaist.ac.jp

## Abstract

An ultra-thin silicon nitride ( $\text{SiN}_x$ ) layer formed by catalytic chemical vapor deposition (Cat-CVD) is used to replace the Si dioxide ( $\text{SiO}_2$ ) layer of a tunnel oxide passivated contact (TOPCon) solar cell. The passivation quality of c-Si with a stack of the ultra-thin  $\text{SiN}_x$  and n-type amorphous or microcrystalline Si, also formed by Cat-CVD, is significantly improved by annealing at 850 °C, probably due to the formation of a back surface field (BSF) layer. Cat-CVD  $\text{SiN}_x$  with thicknesses of up to 2.5 nm can have sufficient tunneling conduction, which is much larger than the upper limit thickness of a tunneling  $\text{SiO}_2$  of ~2 nm. The ultrathin  $\text{SiN}_x$  having functions of surface passivation and carrier tunneling and the unification of the formation method for the tunneling and conductive layers will lead to the realization of tunnel nitride passivated contact solar cells.

## 1. Introduction

Recently, crystalline silicon (c-Si) solar cells with a full-area carrier selective contact passivated by a tunneling silicon dioxide ( $\text{SiO}_2$ ) layer on rear side have been proposed, which are referred to as tunnel oxide passivated contact (TOPCon) solar cells [1]. To form the passivated contacts, there are three process steps: (i) formation of an ultra-thin  $\text{SiO}_2$  layer, usually performed by wet-chemical oxidation or thermal oxidation, (ii) deposition of a doped Si layer, in which plasma-enhanced chemical vapor deposition (PECVD) and low-pressure CVD (LPCVD) are commonly used, (iii) an indispensable “activation” for this contact by high-temperature annealing. The fabrication process of TOPCon solar cells can be simplified by unifying the method of thin-film formation in the steps (i) and (ii).

Catalytic CVD (Cat-CVD) is one of the best ways for the formation of tunneling and doped Si films, because of its advantages such as higher material efficiency and plasma-damage-less nature. Si nitride ( $\text{SiN}_x$ ) films formed by Cat-CVD is known to have excellent passivation ability on c-Si surfaces [2], and the replacement of  $\text{SiO}_2$  in TOPCon solar cells to ultra-thin Cat-CVD  $\text{SiN}_x$  will lead to the simplification of the fabrication process [3]. In this study, we investigated the passivation quality and tunneling conductivity of Cat-CVD doped Si/ $\text{SiN}_x$  passivated contact. We mainly focused on the influences of  $\text{SiN}_x$  thickness and the phosphorus doping concentration of the doped Si layer.

## 2. Experimental details

We used mirror-polished n-type 1–5  $\Omega\text{cm}$  (100)-oriented floating-zone (FZ) Si wafers with a thickness of 280  $\mu\text{m}$ ,

which were cleaved into 20×20 mm<sup>2</sup> pieces. After RCA cleaning, the wafers were carried into a Cat-CVD chamber immediately. The wafers were pre-heated up to a setting temperature ( $T_{\text{sub}}$ ) in hydrogen atmosphere, and then  $\text{SiN}_x$  layers were deposited on both sides of the wafers.  $\text{SiN}_x$  layers with various thicknesses in a range of 1.8–4.2 nm were deposited to investigate the dependence of  $\text{SiN}_x$ . Subsequently, n-type amorphous Si (a-Si) or microcrystalline Si ( $\mu\text{c-Si}$ ) films with a thickness of ~25 nm were deposited on both sides of  $\text{SiN}_x$  layers.  $T_{\text{sub}}$ , catalyzer temperature ( $T_{\text{cat}}$ ), pressure ( $P_g$ ), and gas flow rates for the deposition of  $\text{SiN}_x$ , a-Si, and  $\mu\text{c-Si}$  films are summarized in Table I.

To activate the contact, an annealing was conducted in a furnace at 850 °C for 1 hour in  $\text{N}_2$  atmosphere, by which a-Si or  $\mu\text{c-Si}$  will be crystallized into polycrystalline Si (poly-Si). The effective minority carrier lifetime ( $\tau_{\text{eff}}$ ) of the c-Si samples was measured by microwave photoconductivity decay ( $\mu\text{-PCD}$ ) before and after the annealing. Ti/Ag electrodes with an area of 1 cm<sup>2</sup> were deposited on both sides by vacuum evaporation. The  $I$ - $V$  characteristics of the samples were measured using a semiconductor parameter analyzer, from which carrier tunneling through thin  $\text{SiN}_x$  was evaluated. We also separately prepared the samples for the evaluation of contact resistance with electrodes on one side of the samples with various intervals. In the latter structure,  $\mu\text{c-Si}$  with a  $\text{PH}_3$  flow rate of 100 sccm was chosen as a doped Si layer.

Table I Deposition conditions for  $\text{SiN}_x$ , a-Si, and  $\mu\text{c-Si}$  films.

	$T_{\text{sub}}$ (°C)	$T_{\text{cat}}$ (°C)	$P_g$ (Pa)	Gas flow rate (sccm)			
				$\text{SiH}_4$	$\text{NH}_3$	$\text{H}_2$	$\text{PH}_3$ (2.25%)
$\text{SiN}_x$	200	1800	1.0	3	50	40	-
a-Si	250	1800	1.0	20	-	-	10–300
$\mu\text{c-Si}$	250	1800	1.0	10	-	10	10–200

## 3. Results and discussion

Fig. 1 shows the  $J$ - $V$  curves of the samples with  $\text{SiN}_x$  with various thickness. The samples with  $\text{SiN}_x$  with thicknesses of 1.8 and 2.5 nm show a linear  $J$ - $V$  characteristics and have resistances of 0.11 and 0.18  $\Omega\text{cm}^2$ , respectively. This clearly demonstrates that  $\text{SiN}_x$  with a thickness of up to 2.5 nm has a sufficient carrier tunneling ability. As the thickness of  $\text{SiN}_x$  increases, the  $J$ - $V$  curves become nonlinear and their slopes near the origin of coordinates decreases severely, which indicates the attenuation of the tunnel effect. The maximum thickness of  $\text{SiN}_x$  for carrier tunneling is ~2.5 nm, which is much larger than that of  $\text{SiO}_2$  of 2.0 nm [5]. According to the theoretical analysis, the most possible reason is that  $\text{SiN}_x$  has a lower band offset for tunneling electrons.

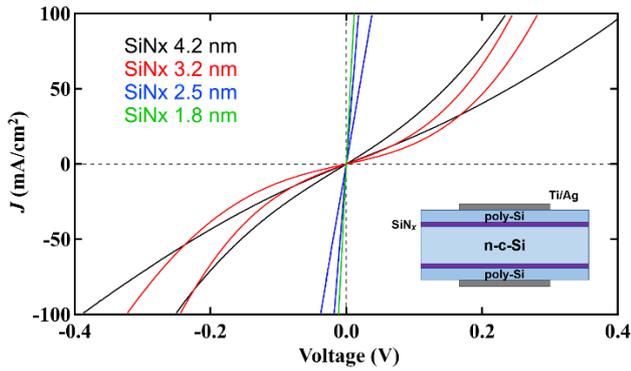


Fig. 1  $J$ - $V$  curves of samples with  $\text{SiN}_x$  with various thickness.

Fig. 2 shows  $\tau_{\text{eff}}$  as a function of  $\text{SiN}_x$  thickness. Annealing at 850 °C leads to a large increase in  $\tau_{\text{eff}}$ , while the increase of  $\tau_{\text{eff}}$  is weakened by the increase in  $\text{SiN}_x$  thickness. The high-temperature annealing may result in the diffusion of phosphorus (P) atoms from doped Si into c-Si and a heavily doped layer is formed on the c-Si surface. This may act as a back surface field (BSF) layer and contribute to an improvement in passivation quality. One possible reason for a decrease in  $\tau_{\text{eff}}$  is that thicker  $\text{SiN}_x$  layer prevents the diffusion of P atoms into c-Si more severely and the field-effect passivation becomes weaker.

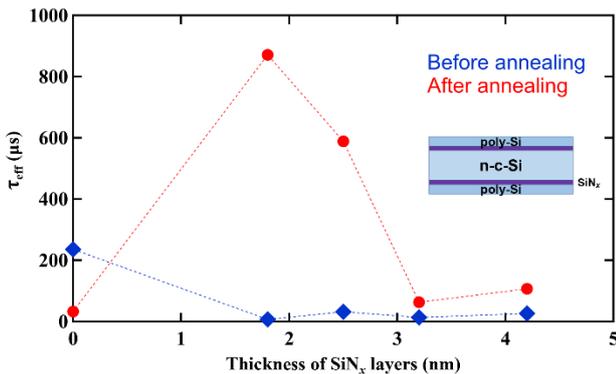


Fig. 2  $\tau_{\text{eff}}$  as a function of  $\text{SiN}_x$  layer before and after annealing.

Fig. 3 shows  $\tau_{\text{eff}}$  after annealing as a function of the He-diluted  $\text{PH}_3$  flow rate during the deposition of a-Si or  $\mu\text{c-Si}$ .  $\tau_{\text{eff}}$  and  $\text{PH}_3$  flow rate, i.e. P concentration, clearly show a strong positive correlation at a low  $\text{PH}_3$  flow rate.  $\tau_{\text{eff}}$  then gradually decreases with further increase in  $\text{PH}_3$  flow rate.

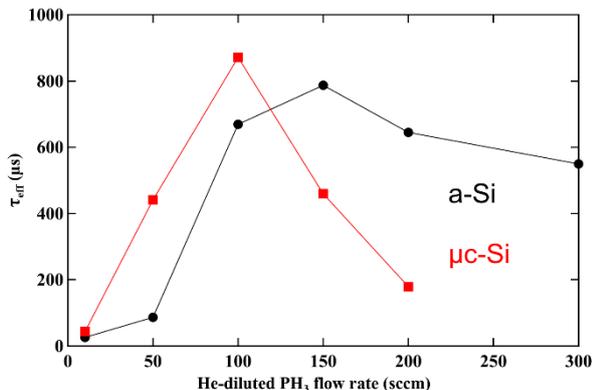


Fig. 3  $\tau_{\text{eff}}$  after annealing as a function of He-diluted  $\text{PH}_3$  flow rate.

The increase in  $\tau_{\text{eff}}$  can be easily understood as the formation of stronger BSF layer by more P diffusion. The successive reduction in  $\tau_{\text{eff}}$  might be related to the worse quality of doped Si films themselves induced by a reduction of the partial pressure of  $\text{SiH}_4$  during the deposition of the Si films. The maximum  $\tau_{\text{eff}}$  of  $\sim 0.8$  ms obtained in this experiment indicates high passivation quality of the n-type Si/ultra-thin  $\text{SiN}_x$  stacks.

Fig. 4 shows the measured resistance of the samples for the contact resistance measurement as a function of electrode interval. The resistivity of the Si substrate ( $\rho_{\text{sub}}$ ) calculated from the slope is 1.03  $\Omega\text{cm}$ , which is lower than that of untreated Si wafer of 3–4  $\Omega\text{cm}$  measured on a quasi-steady-state photoconductance (QSSPC) equipment. One possible reason for the lowering of  $\rho_{\text{sub}}$  is the contribution of lateral conduction in a BSF layer. The contact resistance estimated from the intercept is 0.014  $\Omega\text{cm}^2$ . This value is low enough to for the utilization of c-Si solar cells. The stacked films consisting of n-type Si and ultra-thin  $\text{SiN}_x$ , both prepared by Cat-CVD, can thus combine high passivation ability and sufficient tunneling conductivity, and will be utilized to the passivated contact of c-Si solar cells.

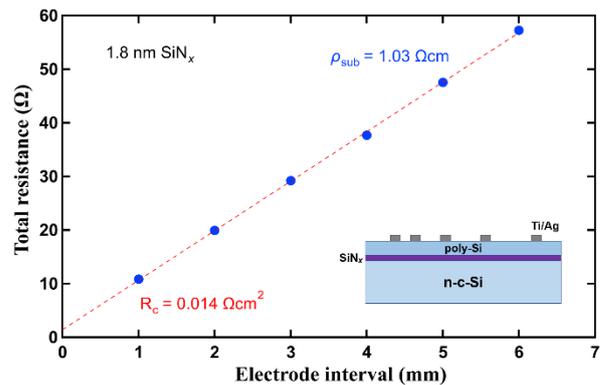


Fig. 4 Measured resistance as a function of electrode interval. The schematic of the sample structure is also shown in the inset. The dashed line indicates the result of linear fitting.

#### 4. Conclusion

We realized ultra-thin Cat-CVD  $\text{SiN}_x$  with functions of high surface passivation ability and sufficient carrier tunneling. The maximum thickness of  $\text{SiN}_x$  for carrier tunneling is found to be as thick as  $\sim 2.5$  nm. The dependence of P concentration on  $\tau_{\text{eff}}$  indicates that the formation of BSF layer may be the primary reason for an improvement in the passivation quality. Contact resistance is estimated to be 0.014  $\Omega\text{cm}^2$ , which is low enough to be used for the passivated contact of c-Si cells, instead of conventional TOPCon structures with  $\text{SiO}_2$ .

#### References

- [1] F. Feldmann *et al.*, Sol. Energy Mater. Sol. Cells **120** (2014) 270.
- [2] T. T. Cham *et al.* Jpn. J. Appl. Phys. **53** (2014) 022301.
- [3] H. Song *et al.*, Jpn. J. Appl. Phys. **57**, (2018), 08RB03.
- [4] S. Mitra *et al.*, IEEE Trans. Electron. Dev. **66** (2019) 1368.
- [5] J. Shewchun *et al.*, J. Appl. Phys. **48** (1977) 765.